



# The frequency of snowline-region planets from four years of OGLE–MOA–Wise second-generation microlensing

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## ABSTRACT

We present a statistical analysis of the first four seasons from a ‘second-generation’ microlensing survey for extrasolar planets, consisting of near-continuous time coverage of 8 deg<sup>2</sup> of the Galactic bulge by the Optical Gravitational Lens Experiment (OGLE), Microlensing Observations in Astrophysics (MOA), and Wise microlensing surveys. During this period, 224 microlensing events were observed by all three groups. Over 12 per cent of the events showed a deviation from single-lens microlensing, and for  $\sim$ one-third of those the anomaly is likely caused by a planetary companion. For each of the 224 events, we have performed numerical ray-tracing simulations to calculate the detection efficiency of possible companions as a function of companion-to-host mass ratio and separation. Accounting for the detection efficiency, we find that  $55^{+34}_{-22}$  per cent of microlensed stars host a snowline planet. Moreover, we find that Neptune-mass planets are  $\sim$ 10 times more common than Jupiter-mass planets. The companion-to-host mass-ratio distribution shows a deficit at  $q \sim 10^{-2}$ , separating the distribution into two companion populations, analogous to the stellar-companion and planet populations, seen in radial-velocity surveys around solar-like stars. Our survey, however, which probes mainly lower mass stars, suggests a minimum in the distribution in the super-Jupiter mass range, and a relatively high occurrence of brown-dwarf companions.

**Key words:** gravitational lensing: micro – surveys – binaries: general – planetary systems – Galaxy: stellar content.

## 1 INTRODUCTION

Over the last 20 years, our knowledge about planetary systems has increased dramatically, from one example with eight planets (our own Solar system) to over 1100 planetary systems hosting more than 1800 planets.<sup>1</sup> The various planet-detection techniques are sensitive to different regions of the planetary parameter space, but the overall emerging picture from combining their results is that

planetary systems are very common. The occurrence rate found by radial velocity and transits surveys, which have detected the majority of the known exoplanets, and are most sensitive to close-in planets, shows that, on average, about half of Sun-like stars host small planets (Earth to Neptune) within 1 AU, 10 per cent have giant planets within a few AU, and about 10 per cent among those giants are within a few hundredths of an AU (Winn & Fabrycky 2015). Moreover, systems with multiple close-in planets are not rare. The existence of close-in massive planets (‘hot Jupiters’) requires a migration mechanism for those planets from beyond the ‘snowline’, where they likely form according to planet formation scenarios (Ida & Lin 2005).

Complementing these discovery methods, gravitational microlensing, while having detected only tens of planets, offers a

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<sup>1</sup> <http://exoplanet.eu>

number of unique advantages. Microlensing is sensitive to planets at projected host separations of 1–10 AU, orbiting stars throughout a large volume of the Galaxy, among both the bulge and the disc populations (Gaudi 2012). Initial statistical analyses of microlensing surveys have suggested that, on average, every star in the Galaxy hosts at least one snowline planet (Cassan et al. 2012), and that the majority of them might be multiple-planet systems (Han et al. 2013). Microlensing surveys are also the only technique sensitive, in principle, to unbound planets. Sumi et al. (2011) argue that such free-floating planets are at least as common as main-sequence stars. However, due to the current small total number of microlensing-detected planets and the difficulties in turning these numbers into a statistical occurrence rate, the abundance of Solar-like planetary systems is still poorly known.

The traditional ‘first-generation’ strategy for exoplanet searches by the microlensing community has been to focus on specific microlensing events discovered by the Optical Gravitational Lens Experiment (OGLE; Udalski 2009) and the Microlensing Observations in Astrophysics (MOA; Sumi et al. 2003) microlensing surveys. These events have generally been bright and of high magnification. Following alerts announced by the surveys, events that promised to be of high magnification ( $A \gtrsim 100$ ) were monitored quasi-continuously by global follow-up networks, in order to search for planetary anomalies. The light curve of a bright and highly magnified event has a high signal-to-noise ratio (SNR), permitting observations with small (and even amateur) telescopes, which are a significant part of the follow-up networks. The motivation for this strategy was that high-magnification events have the highest sensitivity to planets, since the source star image is distorted into a nearly full Einstein ring, and hence a planet anywhere in the Einstein ring vicinity will perturb the light curve and will be detected (Gould & Loeb 1992; Griest & Safizadeh 1998). Although this strategy was important for discovering the first microlensing planets, it also led to the main limitations of the first generation: (1) high-magnification events are rare ( $\sim 1$  per cent) and hence overall a small number of microlensing planets have been discovered with this approach; (2) the alert and follow-up process involves a complex social decision and communication process, sometimes after an event has already hinted at the possibility of an anomaly. This process is difficult to account for in a statistical analysis.

Exploiting the full planet discovery potential of microlensing requires a ‘second-generation’ (hereafter – genII) microlensing survey (Gaudi et al. 2009), in which a large fraction of all ongoing microlensing events towards the Galactic bulge – of both low and high magnification – are monitored continuously with a cadence that is high enough to detect planetary perturbations in the light curves. Apart from the larger number of expected planet discoveries, the ‘controlled experiment’ nature of a genII survey allows for safer statistical inferences on the planet population. The requirements are straightforward – a network of 1m-class telescopes situated so as to allow 24 h coverage of the bulge, equipped with degree-scale imagers.

There have been three main works to constrain snowline planet properties and frequencies based on first-generation surveys. Sumi et al. (2010) used the first 10 microlensing-discovered planets to constrain the mass-ratio distribution of snowline region exoplanets and their hosts, and concluded that Neptune-mass planets are at least three times as common as Jupiters in this region. Gould et al. (2010) analysed six planets discovered in high-magnification microlensing events that were intensively followed-up by the  $\mu$ FUN network (Gould 2008), and found that  $\sim$ one-third of microlensed stars host a snowline planet in the planet-to-star mass-ratio interval  $-4.5 < \log q$

$< -2$ , corresponding to the range of ice to gas giants. Moreover, since one of the systems had two planets with the same mass and distance ratios as Jupiter and Saturn to the Sun, they argued that the frequency of Solar-like systems around lens stars is about one-sixth. Gould et al. (2010) argued that the chaotic nature of the ‘alert and follow-up’ process results, in the end, in a randomization that minimizes selection effects and although it cannot be claimed to be a controlled experiment, it gives a high level of completeness. Finally, Cassan et al. (2012) combined the previous results and three additional planets that were discovered by the PLANET follow-up network (Albrow et al. 1998), and deduced a frequency of more than unity for snowline-region planets.

In 2011 June and July, we carried out a 6-week pilot of a genII experiment. In 2011, we initiated the first full genII microlensing experiment, involving OGLE, MOA and the Wise Observatory microlensing group (Shvartzvald & Maoz 2012). During four observing seasons, we have monitored all events in a specific, 8 deg<sup>2</sup>, field towards the Galactic bulge. In this paper, we present a statistical analysis of these first four observing seasons (2011–2014) of our genII survey. The sample includes 224 microlensing events, monitored by all three groups. Over 12 per cent of the events show clear anomalies in their light curves indicating a companion to the lens star. To assess the nature of the companions we perform a coarse grid-based search of the microlensing model parameter space for each event. Comparing the actual number of light curves having planetary anomalies to the total number of events, after accounting for the detection efficiency of each event, allows us to estimate, for the first time, the frequency of planetary systems from a controlled microlensing experiment.

The paper is arranged as follows. In Section 2, we describe the observations and reductions of our sample of lensing events. In Section 3, we model each light curve in our sample with a single-lens model, describe our anomaly detection filter, and present the anomalous events among our sample. In Section 4, we derive the frequency and the mass-ratio distribution of snowline planets implied by our survey, and discuss the results in the context of current knowledge in Section 5.

## 2 OBSERVATIONS

### 2.1 Observational data and reduction

Our genII survey network is a collaboration between three groups: OGLE, MOA, and Wise. The OGLE and MOA groups regularly monitor a large region of the Galactic bulge, and routinely identify and monitor microlensing events. The Wise group monitors a field of 8 deg<sup>2</sup>, within the observational footprints of both OGLE and MOA, having the highest event rates based on previous years’ observations (see Shvartzvald & Maoz 2012).

The OGLE group operates the 1.3 m Warsaw University telescope at Las Campanas Observatory in Chile, with the 1.4 deg<sup>2</sup> field of view OGLE-IV camera. The genII survey’s 8 deg<sup>2</sup> footprint includes three extremely high-cadence OGLE-IV fields observed every  $\sim 15$  min and five high-cadence OGLE-IV fields observed once per  $\sim 45$  min. Standard OGLE observations are through an *I* filter. OGLE photometry used here was extracted by their standard difference image analysis (DIA) procedure (Udalski 2003). These OGLE data are the preliminary measurements posted on the Early Warning System website,<sup>2</sup> which are not suitable for detailed

<sup>2</sup> <http://ogle.astrouw.edu.pl/ogle4/ews/ews.html>

modelling of the events, but sufficient for the anomaly-detection search and rough modelling that we perform here.

The MOA group operates the 1.8 m MOA-II telescope in New Zealand with the MOA-cam3 camera, with a 2.2 deg<sup>2</sup> field of view (Sako et al. 2008). MOA-II monitors six extremely high-cadence fields (every ~15 min) which cover most of the genII 8 deg<sup>2</sup> field, and an additional six high-cadence fields (every ~45 min) which complement the network’s footprint. MOA uses a wide *R/I* filter (‘MOA-Red’). The MOA data were reduced using their routine DIA procedure (Bond et al. 2001).

The Wise group’s main setup is the 1m telescope at Wise Observatory in Israel equipped with the Large Area Imager for the Wise Observatory (LAIWO) camera, with a 1 deg<sup>2</sup> field of view. The cadence for each of the eight Wise pointings (which define the survey footprint) is ~30 min. Observations are in the *I* band. In the first half of the 2012 season, the LAIWO camera had readout electronics problems and the Wise group used, as an alternative, the Wise C18 0.46 m telescope with an *I* filter, with a 1 deg<sup>2</sup> field of view. The cadence for the eight C18 fields (which overlap with, but differ from the eight LAIWO fields) was ~1 h. All Wise data were reduced using the PYSIS DIA software (Albrow et al. 2009).

## 2.2 Sample of lensing events

The sample of microlensing events analysed here consists of 224 events from the 2011–2014 bulge seasons, observed by all three groups, and with each group having data near the peak of the event. Without the last criterion, events with a long time-scale that peak before the beginning of a season might enter the sample, while such events with a short time-scale would not. Being the northernmost node of the network, with the shortest bulge observing season, Wise sets the time limits for the genII survey, from mid-April to mid-September of each season.

With the above criteria, and considering the limiting magnitudes of the various groups (see definition in Section 3.1), the initial sample has a relatively small number of events (230). After identifying anomalous events in our sample (see Section 3.3), we noted six events that passed the limiting magnitude criteria only because of the anomaly itself (due to a caustic crossing), and otherwise would not meet the criteria. Including these events would introduce a bias towards anomalous events, and therefore we exclude them from the sample. Table A1 lists the final sample of 224 events. For completeness, the excluded six events are listed separately at the end of the table.

## 3 COMPANION DETECTIONS

### 3.1 Single-lens modelling of all events

As a first stage in our analysis, we model each of the 224 events with a single point-lens model, including the possibility of parallax and finite source effects. Such a model is necessary in order to subsequently detect those events that have lensing anomalies (Section 3.3), as well as to gauge the planet detection efficiency of each event (Section 3.2.2). Both of these elements of the analysis then enter our planet frequency calculation (Section 4).

The characteristic angular scale of a lensing phenomenon is the angular Einstein radius,

$$\theta_E^2 = \kappa M \pi_{\text{rel}}, \quad \kappa = \frac{4G}{c^2 \text{au}} \simeq 8.14 \frac{\text{mas}}{M_\odot}, \quad (1)$$

where  $M$  is the mass of the lens star, and the relative parallax is  $\pi_{\text{rel}} = \text{au}(D_L^{-1} - D_S^{-1})$ , where  $D_S$  and  $D_L$  are the distances to the source and lens stars, respectively. A standard point-source, point-lens, microlensing model requires three ‘Paczynski parameters’:  $t_E$  – the event time-scale, over which the source crosses one angular Einstein radius;  $u_0$  – the impact parameter, the minimal separation between the source star and lens star in angular Einstein radius units; and  $t_0$  – the time of minimal separation between the source star and the lens star. There are two higher order effects, which can change the light curve even in the absence of a companion to the lens. First, the finite source size needs to be taken into account when the scale in the source plane over which the magnification varies is of the order of the source size. This effect is parametrized by the ratio between the angular size of the source,  $\theta_*$ , and the angular Einstein radius,  $\rho = \theta_*/\theta_E$ . Secondly, the orbital microlens parallax,  $\pi_E$ , changes the relative proper motion between the source star and the lens during an event, due to the Earth’s orbit. The effect is parametrized by the size of the angular Einstein radius, projected on to the observer plane in astronomical units (AU), and can be represented by:  $\pi_E = (\pi_{\text{rel}}/\theta_E)(\mu/\mu)$ , where  $\mu$  is the relative proper motion between the source and the lens. The microlens parallax is divided into a north and east component with respect to the Galactic coordinate system,  $\pi_{E,N}$  and  $\pi_{E,E}$ . The orbital parallax involves a well-known  $u_0 \leftrightarrow -u_0$  ecliptic degeneracy for sources near the ecliptic (Jiang et al. 2004; Poindexter et al. 2005), and we explore both options in our modelling.

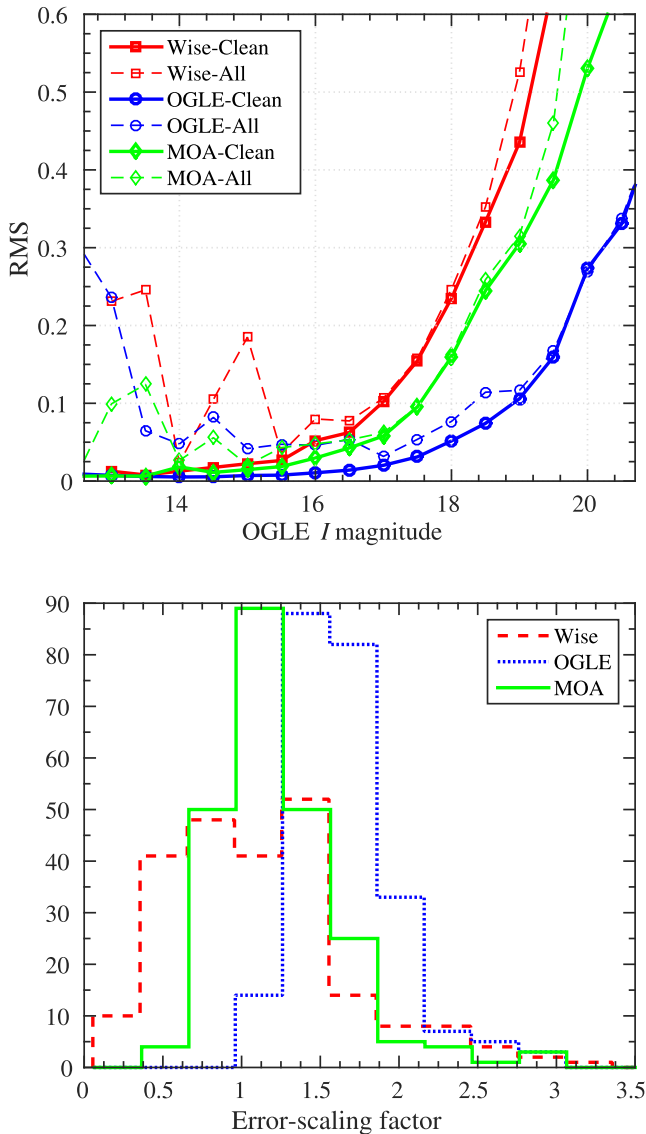
During a microlensing event, an additional constant flux from stars along the line of sight to the source, that do not get magnified due to the lens (in particular, the lens star itself), contributes to the light curve. Thus the total flux,  $f(t)$ , is the sum of the source flux,  $f_s$ , times the magnification,  $A(t)$ , and the background flux,  $f_b$ . Due to the different filters and seeing conditions of each survey group, this is calculated separately for each data set,  $i$ ,

$$f_i(t) = f_{s,i}A(t) + f_{b,i}. \quad (2)$$

When exploring the lens-model parameter space, we limit the ranges of some of the parameters, to avoid biases and unphysical solutions. First, we include only events with  $|u_0| < 1$ , i.e. with peak magnifications  $A > 1.34$ . In addition, we limit the crossing time-scale to values  $t_E < 500$  d. Finally, we limit each component of  $\pi_E$  to be less than 1.5. Events that exceed these limits on  $t_E$  and  $\pi_E$  can, in principle, occur for extreme configurations. However, their probability is negligible given the number of events in our sample.

The MOA and Wise fluxes were aligned to the OGLE *I*-band magnitude scale, and inter-calibrated to the single-lens microlensing model. In order to obtain the different photometric precision for each group, we construct an observed distribution of the residuals from the model for the entire data set (from all events), and calculate its root mean square (rms) as a function of OGLE *I*-band magnitude. This accounts for differences in telescope aperture, the different seeing conditions at each site, and overall systematics in the reduction pipelines. Fig. 1 shows the rms of the residuals for each group.

Yee et al. (2012) have studied the effect of faint sources on microlens parameters for the lensing event MOA-2011-BLG-293, and found that when the measured flux errors are comparable to, or larger than the source flux, particularly near baseline, this leads to biases in the measured Einstein time-scale. Based on their conclusions and by analysing the precision at each site, found above, we exclude from each light curve in our sample data points for which the observed residual distribution corresponding to the underlying event’s *I* magnitude (as suggested by the single lens model) has an



**Figure 1.** Top: root mean square (RMS) of residuals of best fits of single-lens microlensing models to all observed light curves, including all data (dashed line) and after excluding outliers and anomalies (solid line), for each group as a function of OGLE  $I$ -band magnitude at a given point in the light curve (OGLE – blue circles, MOA – green diamonds, Wise – red squares). The differences between groups reflect both the instrumental and site quality (collecting area, seeing conditions, CCD sensitivity) and the DIA pipeline for extracting the photometry. Bottom: error-scaling factor distributions for the 224 events, and for each group. Most of the pipeline-reported errors underestimate the real uncertainties. The tails shown by each distribution at large factors reflect events with long-duration anomalies.

rms) greater than 0.33 mag (SNR $\sim$ 3, below that the microlensing signal cannot be reliably detected). This corresponds to events with Wise points with baseline magnitudes fainter than 18.5 mag, and MOA baseline points fainter than 19.5 mag. For OGLE, we include all points.

Each microlensing event in our sample is modelled with the three Paczyński parameters, the two flux inter-calibration parameters for each data set, the finite source effects, and the orbital parallax (testing for both  $u_0 > 0$  and  $u_0 < 0$ ). We explore this parameter space and find the best model parameters, and their uncertainties, using a Markov Chain Monte Carlo (MCMC) algorithm. Outlier data points

and real microlensing anomalies will obviously draw our solution away from the true one. We therefore use an iterative algorithm for the exclusion of outliers (which naturally also applies to intervals in the light curve that include anomalies). The photometric errors reported by DIA programs are usually underestimates. Therefore, for each data set, we scale these errors to set  $\chi^2$  per degree of freedom (dof) equal unity, without the outliers and anomalies which will bias the distributions, allowing us to estimate the real precision of each data set. Fig. 1 shows the distribution of error scaling factors for each group. The median factors for OGLE, MOA and Wise are 1.74, 1.32 and 1.21, respectively. Finally, we initiate the entire modelling process again without the excluded outliers and with the rescaled errors. Fig. 2 shows an example of the results of our outlier detection and exclusion algorithm, for an event with a strong parallax signal, where the single-lens model fits well the light curve after outlier points are excluded.

Table A1 in Appendix A lists the best single-lens model parameters for each event in our sample, and their uncertainties, which were extracted from the MCMC estimation of the covariance matrix. Microlens parallax is significantly detected ( $\Delta\chi^2 > 100$  with respect to a model without parallax) in 10 per cent of the non-anomalous events, while finite source effects are found in 3 per cent of the non-anomalous events. The values for the microlens parallax and finite source parameters for those events are marked in boldface in the table. We do not mark the values for anomalous events, since the anomaly might affect those parameters, and they can be reliably detected only with a full binary-lens model.

Fig. 3 shows, for the entire sample of events, the distributions of the best-fitting values of the event time-scale,  $t_E$ , and the impact parameter,  $u_0$ . The  $t_E$  distribution is shown fitted with a lognormal function,  $N \propto \exp(-\frac{(\log(t_E) - \mu)^2}{2\sigma^2})$ , with median time-scale of 19.5 d.  $u_0$  should, in principle, follow a uniform probability. However, faint sources that undergo high magnification will be detected, while similar sources undergoing a lower magnification event will be missed, due to the limiting magnitudes of the surveys, introducing a bias towards high-magnification (low  $u_0$ ) events. The  $t_E$  and  $u_0$  distributions of our sample are almost identical to the distributions of those parameters from previous years’ OGLE and MOA microlensing seasons (see Shvartzvald & Maoz 2012). This shows that our sample has no obvious biases, and thus represents a typical population of lenses and sources.

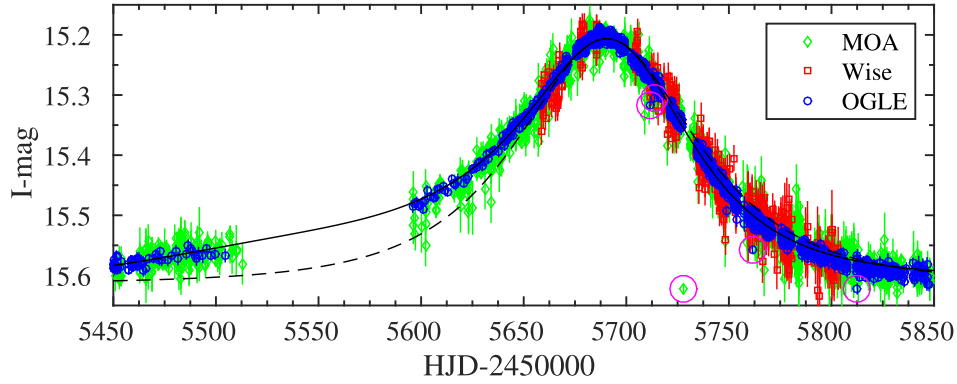
### 3.2 Anomaly detection filter

A companion (or multiple companions) to the lens star can break the symmetry of the microlensing light curve, introducing an anomaly with respect to the single-lens model. In order to identify microlensing anomalies in the light curves in our sample, we follow Shvartzvald & Maoz (2012) and define an anomaly detection filter using a ‘running’  $\chi^2$  estimator,

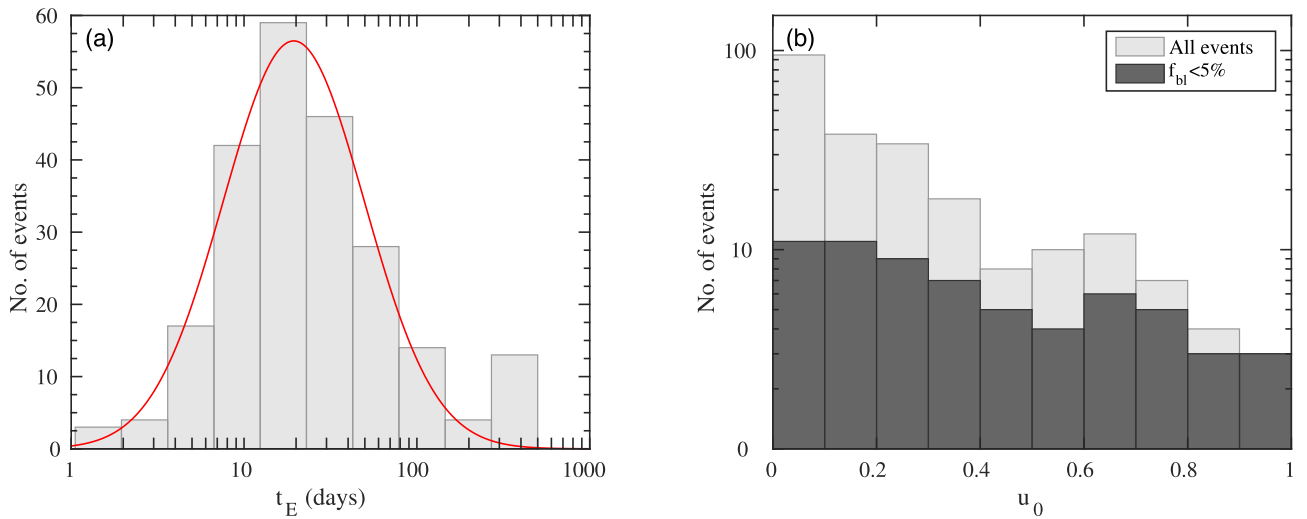
$$\chi_{\text{local}}^2 = \sum_{i=1}^N \frac{(f_i - f_{\text{pl}})^2}{\sigma_i^2}, \quad (3)$$

where  $f_i$  and  $\sigma_i$  are the observed flux and the re-scaled error of each epoch, respectively, and  $f_{\text{pl}}$  is the point-lens model flux at that time. The local  $\chi^2$  is repeatedly calculated by advancing the centre of the filter one observed epoch at a time. Since the relative duration of an anomaly is roughly proportional to the square root of the lens-companion mass ratio, we would like the time interval of each summation to be as short as possible, in order to be sensitive to low-mass planets. However, in order to avoid false-positive detections





**Figure 2.** Example of detection and exclusion of outlier points from the single-lens model in event OGLE-11-0022/MOA-11-025, which has a strong microlens parallax signal (but no detected anomaly from a companion). The detected outlier points are circled. The solid line is the best-fitting microlensing model including parallax, and the dash-dotted line is one without parallax. The light curve is clearly asymmetric, and thus without including parallax it would have been flagged as anomalous.



**Figure 3.** Best-fitting single-lens parameters for the sample of 224 lensing events: (a) distribution of event time-scale,  $t_E$ . The solid curve is a lognormal fit,  $N \propto \exp(-\frac{(\log(t_E) - \mu)^2}{2\sigma^2})$ , with  $\mu = 1.29$  (corresponding to 19.5 d) and  $\sigma = 0.54$ ; (b) distribution of impact parameter,  $u_0$  (light-grey). Due to selection favouring the detection of highly magnified faint events over weakly magnified faint events, there are more events at smaller impact parameters. The dark-grey histogram shows the distribution only for unblended (i.e. relatively bright) events, which are distributed more uniformly, as expected.

in short time-scale events,  $N$  cannot be too small, and we study the sensitivity of our filter to  $N$  below (see Section 3.2.1 for more details). An anomaly is considered detected, if:

- (1)  $\chi^2_{\text{local}}$  divided by the number of summed points,  $N$ , is larger than some threshold value,  $P_{\text{thresh}}$ ;
- (2) at least three consecutive data points have at least a  $3\sigma$  deviation from the point-lens model.

This latter criterion deals with the possibility that photometric outlier points will affect the local  $\chi^2$  test and can lead to false detections.

### 3.2.1 False-positive optimization

In order to optimize the number of false-positive detections (i.e. to reduce their numbers without overly sacrificing true detections), we have studied the sensitivity of our anomaly detection filter to its two parameters ( $N$ ,  $P_{\text{thresh}}$ ). For each of the 224 events in our sample, we use the fitted point-lens parameters to construct a magnification curve. We then sample the light curve at the exact epochs on which the event was observed by each group. With this we account for

all gaps due to weather or technical failures, and for the different sampling cadence by each group and for each field. The theoretical magnification curve is then randomly noised, using the observed residual distribution that we have found for each group (see Section 2.1), and assigned the scaled errors from the real data, to create 500 simulated point-lens light curves for each event. We then search for a false-positive anomaly in the simulated light curve (which has no real anomalies), varying the filter parameters.

We find that the fraction of simulated light curves that trigger the photometric outlier filter (i.e. three consecutive points with  $>3\sigma$ ) is 0.65 per cent. Examining those triggers, we find that 92 per cent of them have  $\chi^2_{\text{local}}/N < P_{\text{thresh}}$ , where  $P_{\text{thresh}}$  is in the range of 7.5–8. The optimal  $P_{\text{thresh}}$  varies from event to event, and depends linearly on the maximum amplification ( $D_{\text{mag}}$ ), which is the difference between the observed baseline magnitude of the event and its peak magnitude, with also a weak dependence on  $N$ . The optimal value of  $N$  is  $N = 40$ . For much smaller  $N$ , the  $\chi^2_{\text{local}}$  can be dominated by a single outlier point, while for much larger values, the anomalous region can be smoothed out and not detected. Using this  $N$  and the  $D_{\text{mag}}$ -dependent value of  $P_{\text{thresh}}$ , 8 per cent of the 0.65 per cent of the

outlier-triggered events trigger an anomaly detection, and thus the resultant false-positive rate is at a low level of 0.05 per cent. These high-threshold criteria help to reduce false-positive detections in the real data, but do not completely avoid them, due to the presence of red noise (i.e. correlated deviant points), a noise that is not included in these randomly noised simulations.

### 3.2.2 Detection efficiency

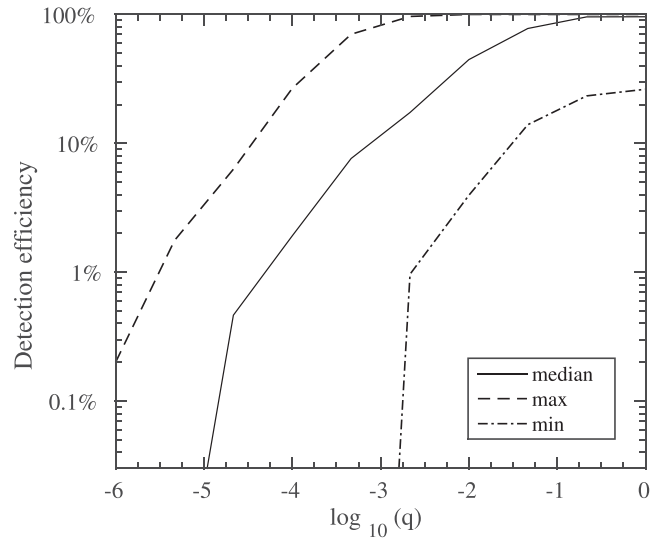
In order to estimate the sensitivity of our microlensing survey to planets and stellar binaries, we have constructed a large sample of simulated microlensing light curves for point-mass lenses with point-mass companions, which were calculated with an adaptive-mesh inverse ray-shooting microlensing light-curve generator (see Shvartzvald & Maoz 2012 for details). Specifically, for each of the 224 events in the real sample, we created a model with the same best-fitting point-lens parameters, as described above, plus a companion to the lens star. The companion introduces three additional parameters:  $q$  – the mass ratio between the companion and the lens star; and the two-dimensional projected position of the companion in the lens plane, relative to the host star position, given by means of a projected angular separation,  $s$ , in units of the angular Einstein radius, and an angle,  $\alpha$ , measured counter-clockwise from the source trajectory in the lens plane. We explore a mass-ratio distribution, uniform in  $\log q$ , in the range  $-6 < \log q < 0$ , a uniform distribution in the scaled projected separation of  $0.3 < s < 3$  (which encompasses the region microlensing is sensitive to), and all possible angles  $\alpha$ . Our working assumption is that the probability for a companion is independent of an event’s single-lens parameter probabilities, which incorporate the host-star properties of mass, proper motion, and distance. For every real event, 3000 simulated light curves, with a variety of companions, were generated.

The simulated magnification curves of each event are then sampled exactly like the real event and randomly noised, as described for the point-lens models used for the false-positive optimization procedure. We then search for anomalies in the simulated light curves with the same procedure and detection criteria that were used for the real data. Explicitly, each light curve is modelled with a point lens and higher order effects (including error-scaling for each observatory), and we search for an anomaly with our detection filter, using the parameters we found for minimizing the false-positive rate.

The ‘lensing zone’ is the range of separations,  $s$ , in which microlensing is sensitive to anomalies caused by a given mass ratio. For stellar binaries (i.e.  $q \geq 10^{-2}$ ), the lensing zone can extend over a factor of a few in  $s$ . For small companions,  $q \leq 10^{-3}$  (likely planets), we find that the range  $0.5 \leq s \leq 2$  covers over 95 per cent of the detections, and therefore, we calculate our detection efficiency for that range. The companion detection efficiency for each event as a function of mass ratio  $q$  was found by marginalizing over  $s$  and  $\alpha$ :

$$\eta(q) = \frac{\iint \Theta(q, s, \alpha) s \, ds \, d\alpha}{\iint s \, ds \, d\alpha}, \quad \Theta(q, s, \alpha) = \begin{cases} 1, & \text{anomalous} \\ 0, & \text{no anomaly.} \end{cases} \quad (4)$$

Fig. 4 shows the results of our detection efficiency simulation. The sensitivity varies by over an order of magnitude among different events, due to the brightness of the event, its magnification and its time-scale, as shown by the 90th and 10th percentile efficiency curves (dashed) in Fig. 4. We find that our experiment is  $\sim$ six times more sensitive to  $q = 10^{-3}$  (corresponding to a  $\sim$ Jupiter/Sun mass ratio) than to  $q = 10^{-4}$  (‘super-Neptunes’).



**Figure 4.** Detection efficiency for a single companion as a function of mass ratio,  $q$ . The event with median efficiency is shown as the solid line. The dashed curves represent events with 10th and 90th percentile sensitivities.

## 3.3 Anomalous events

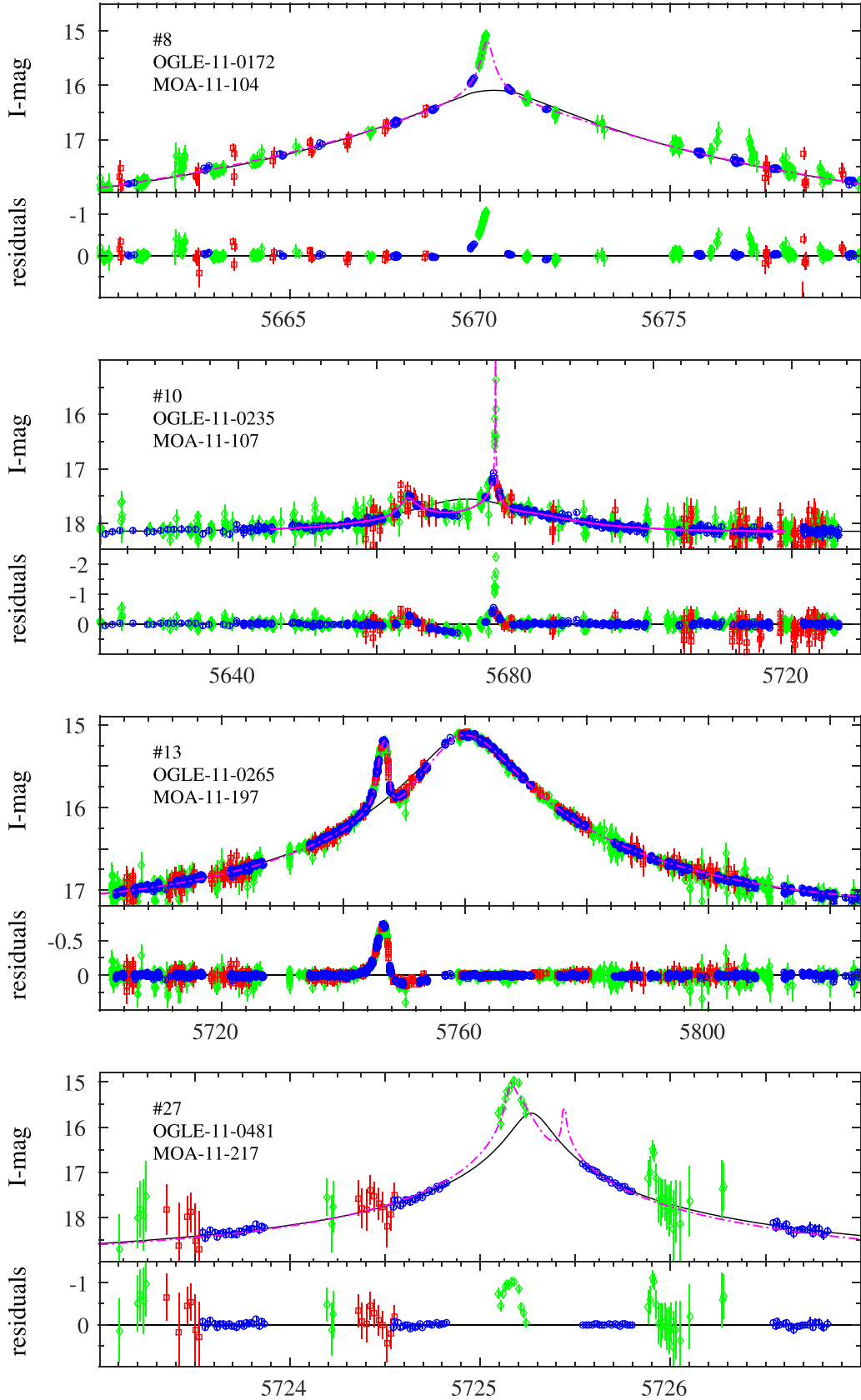
### 3.3.1 Heuristic detections

Before applying to the sample our objective automated anomaly detection filter, we have identified ‘by eye’ 26 anomalous events among the sample in which deviations from the best-fitting point-lens models are clear. Figs 5–11 show the inter-calibrated light curves of these 26 events and their best-fitting single-lens models, obtained after excluding the anomalous regions from the fit (i.e. points detected by the anomaly detection filter). Eight of those anomalous events have already been modelled and published, as follows.

(i) MOA-2011-BLG-293 (Yee et al. 2012) was the first planet discovered by our genII survey. The event was highly magnified, and so was also followed up by several groups. However, Yee et al. (2012) showed that the survey data alone were sufficient to fully constrain the planetary model. Batista et al. (2014) used Keck adaptive optics observations to furthermore constrain the system’s parameters by isolating the light from the lens star, and found that this  $\sim 5M_J$  planet is actually the first microlensing planet discovered in the habitable zone of its host star (a G-type main-sequence star).

(ii) MOA-2011-BLG-322 (Shvartzvald et al. 2014), a moderate-magnification event that did not trigger alerts and follow-up efforts, was the first planetary microlensing event that was detected and analysed based solely on the genII survey data. Using a Bayesian analysis that incorporates a Galactic structure model, Shvartzvald et al. (2014) estimated that it is a  $\sim 12M_J$  planet orbiting an M-type dwarf.

(iii) OGLE-2011-BLG-0265 (Skowron et al. 2015) is a Jupiter-mass planet (with two degenerate solutions of  $\sim 1M_J$  and  $\sim 0.6M_J$ ) orbiting an M dwarf. After the main anomaly was realized from the genII data, the event was alerted. Several follow-up groups monitored the event, allowing a better characterization of the secondary anomaly. This is an example of the problematic, in terms of an eventual statistical interpretation, process of the ‘first-generation’ mode of surveys. The survey data alone were, again, sufficient to fully constrain the planetary model.



**Figure 5.** Inter-calibrated light curves of events with a clear anomaly that we identify by-eye: OGLE – blue circles, MOA – green diamonds, Wise – red squares. Magenta lines are the best-fitting model, either published or from the grid search. The residuals from the point-lens model are shown in the lower panel of each event. The identifying number from Appendix A is marked in the upper-left corner.

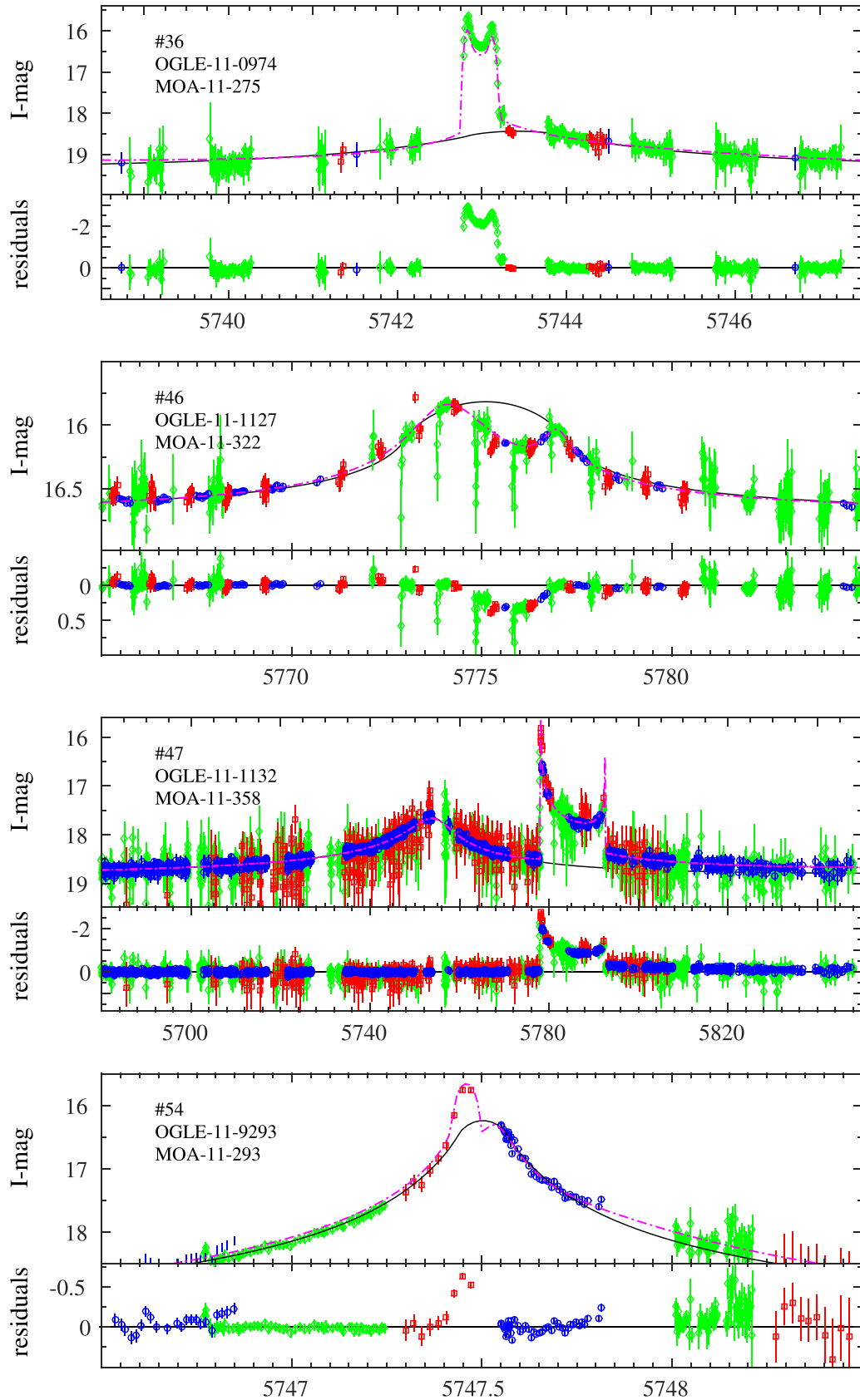


Figure 6. Continued from previous page.



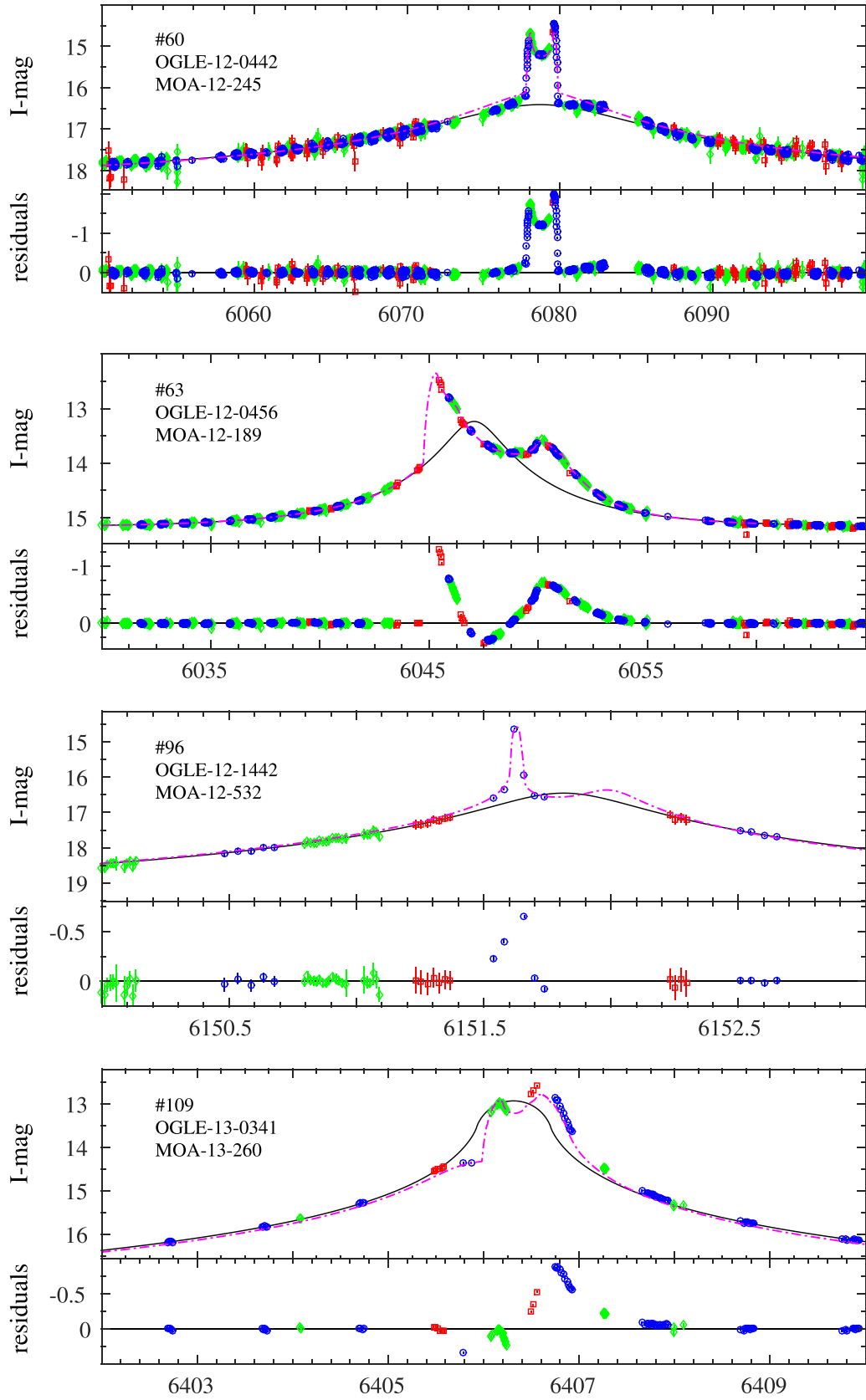


Figure 7. Continued from previous page.

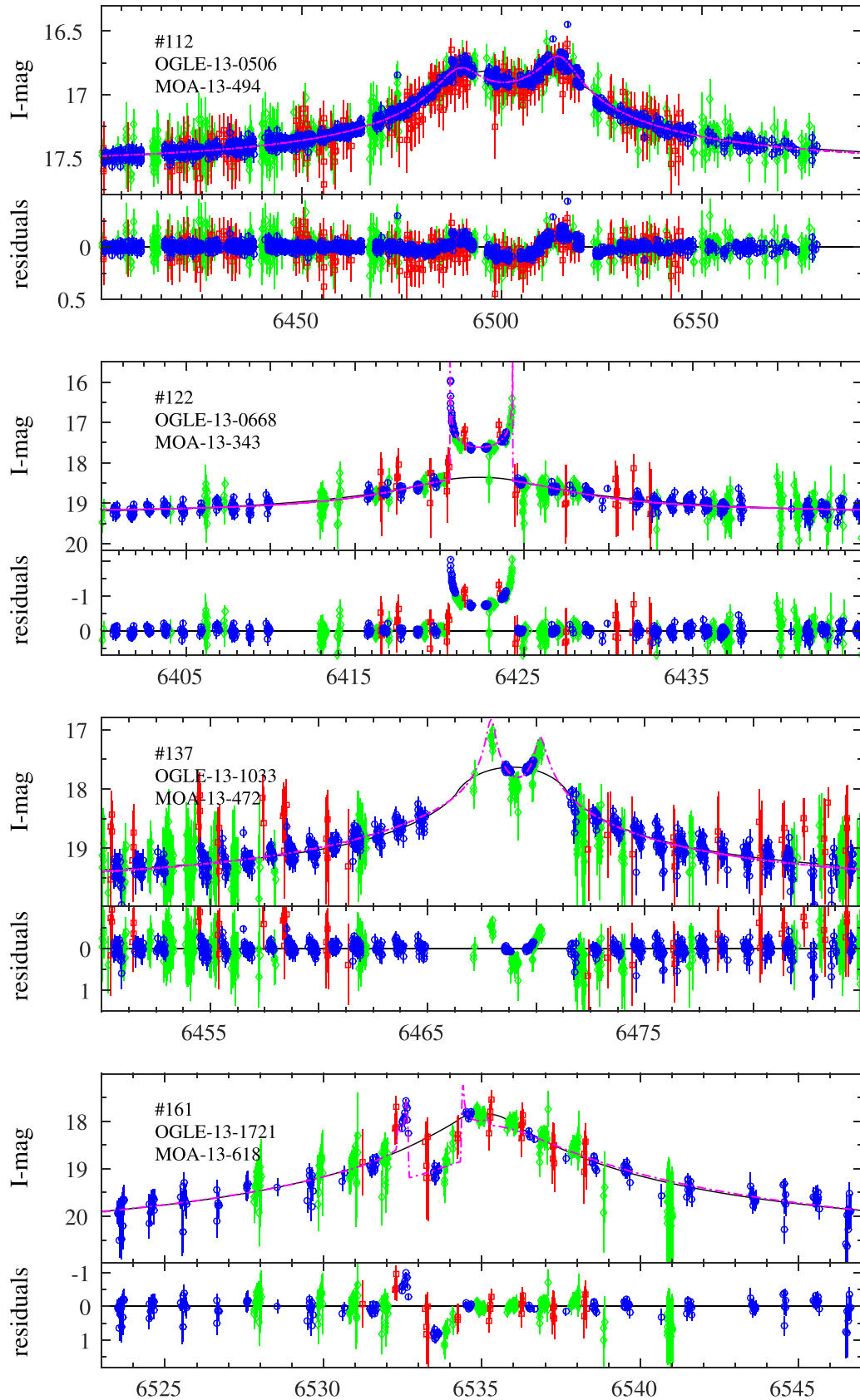


Figure 8. Continued from previous page.

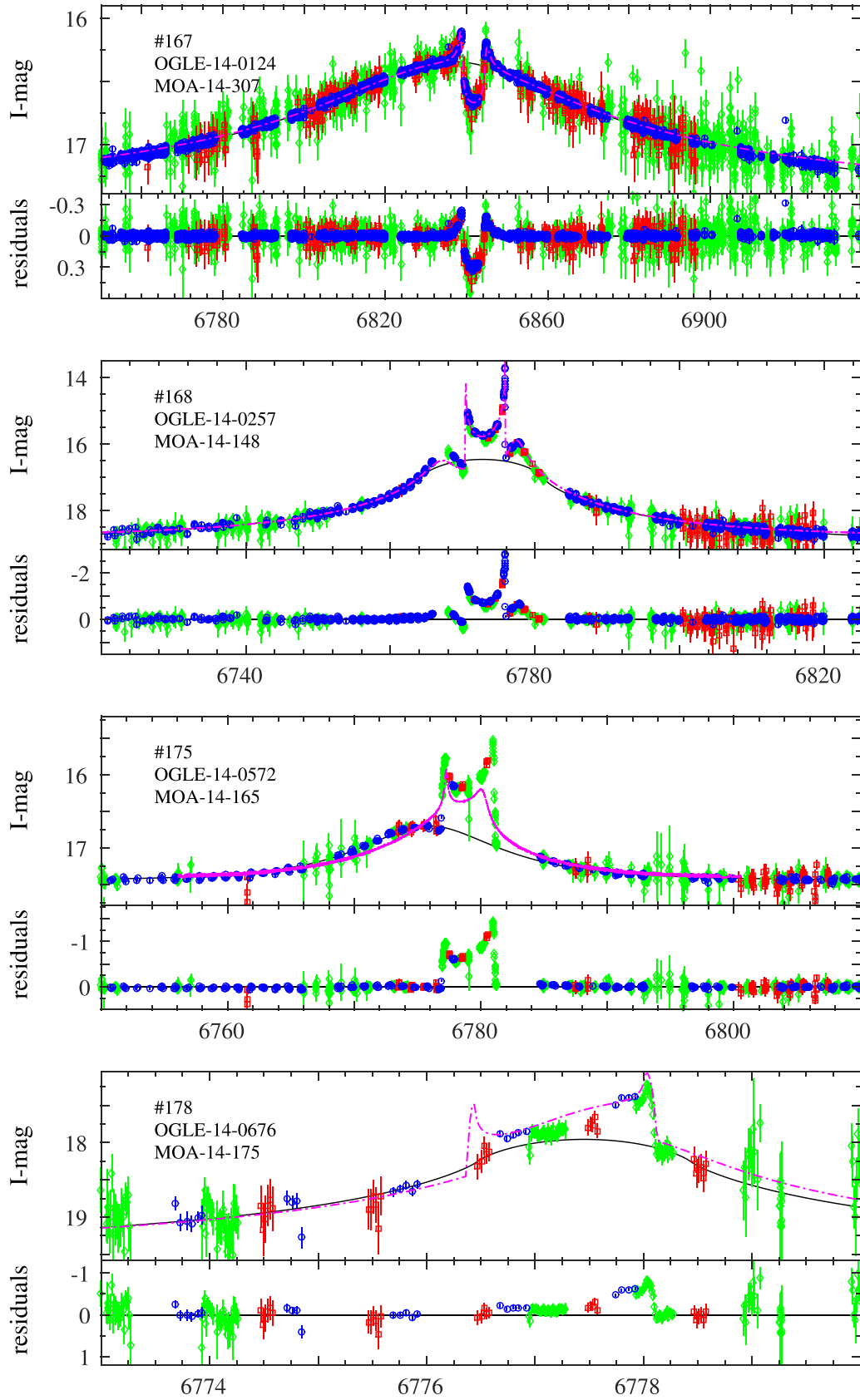


Figure 9. Continued from previous page.

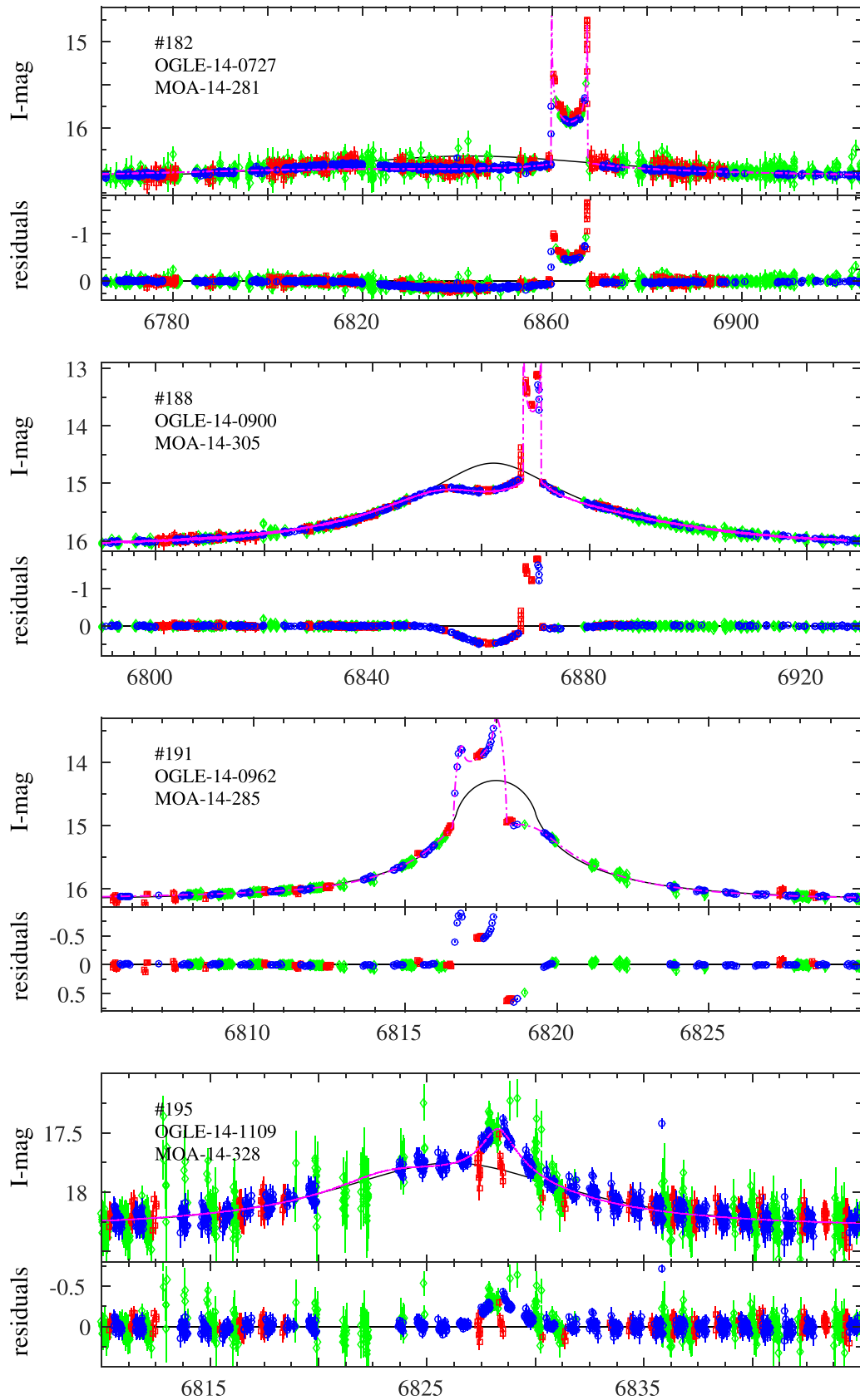


Figure 10. Continued from previous page.

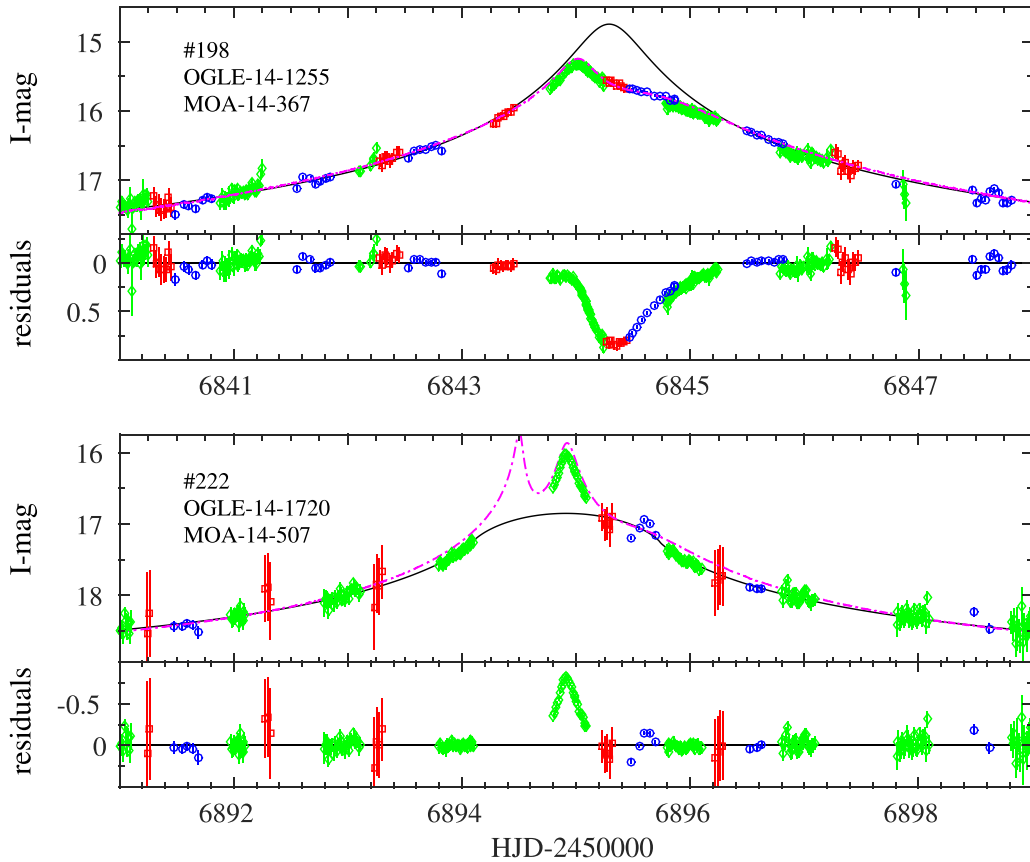


Figure 11. Continued from previous page.

(iv) OGLE-2013-BLG-0341 (Gould et al. 2014) is a  $\sim 2M_{\oplus}$  planet in an  $\sim 1$  AU orbit around one member of a  $\sim 15$  AU separation binary M-dwarf system. The event was alerted and followed up since it was highly magnified, but the genII survey data cover all three anomaly regions. Moreover, the two lower amplitude features, which are the ones that revealed the presence of a planet, were covered only by the genII survey data.

(v) OGLE-2014-BLG-0124 (Udalski et al. 2015) is a  $\sim 0.5M_J$  planet orbiting a K-type dwarf. This was the first microlensing planet co-observed with the *Spitzer* space mission, and used to measure the microlens parallax from simultaneous observations from Earth and space. Udalski et al. (2015) use data only from OGLE and *Spitzer*, but the anomaly and the parallax signal are obvious also from the genII survey data alone.

(vi) The physical parameters of the companions in three additional published events have not been fully determined, but all of them are likely to be binary-star systems. MOA-2011-BLG-104 (Shin et al. 2012) is a binary with a candidate brown-dwarf companion ( $q \sim 0.09$ ). OGLE-2012-BLG-0456 and MOA-2012-BLG-532 (Henderson et al. 2014) have mass ratios of  $q \sim 0.9$  and  $0.5$ , respectively.

We note that the anomalies seen in OGLE-2011-BLG-0481/MOA-2011-BLG-217, which appear only in MOA data, are possibly not real, and if not are an example of the type of false positives that can enter our sample after all. In line with our methodology, we include this event in our sample.

### 3.3.2 Objective detections

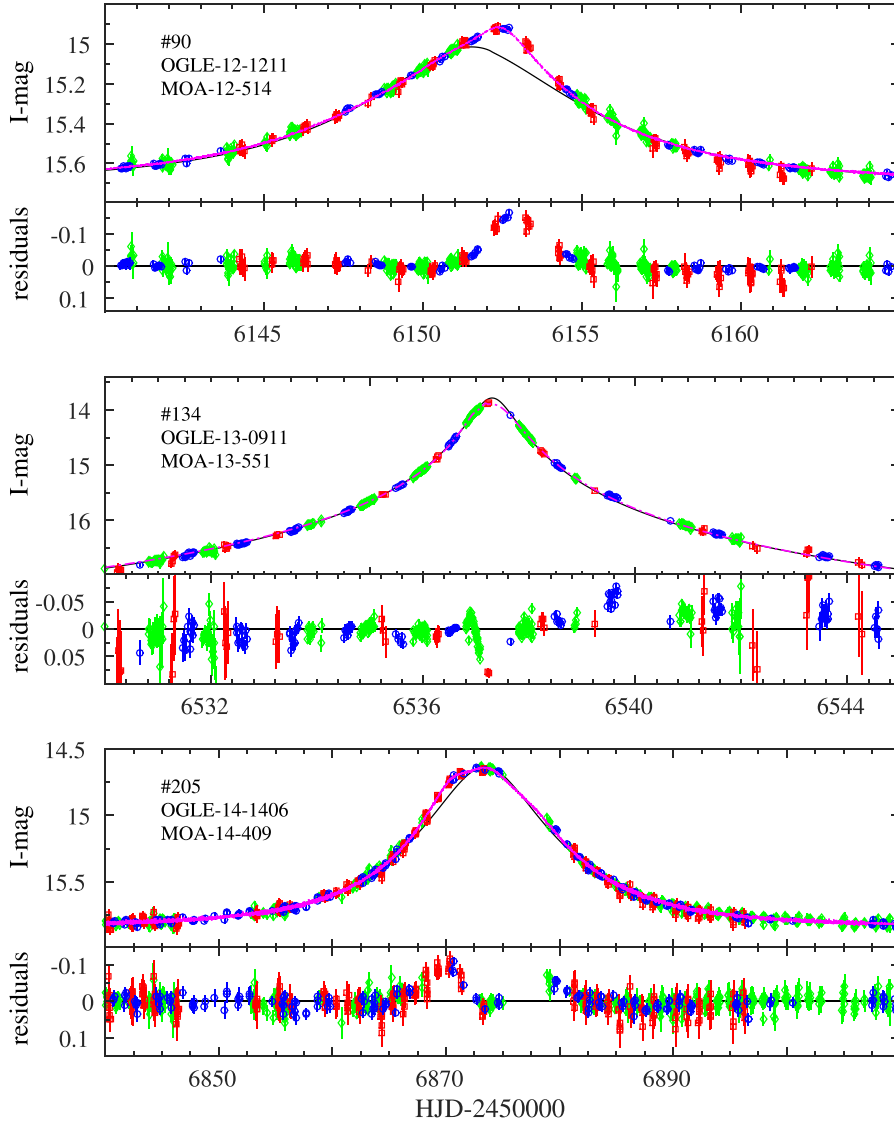
We next apply our automated detection filter to search for anomalies in each event in our sample, with the optimized false-positive filter parameters that we found. The 26 anomalous events, identified by eye in Section 3.3.1 are all easily re-detected by the automated filter. In addition, the automated filter detects three more anomalous events. The light curves of those events, and the residuals from their best-fitting point-lens models, are shown in Fig. 12.

Table 1 summarizes all anomalous events and their estimated mass ratios. The mass ratios are either from the full binary-lens-published models (eight events, see Section 3.3.1) or from a grid search of binary-lens models from a library we constructed, similar to the one used in Section 3.2.2 for the detection efficiency estimates, with 36 000 combination of  $s$ ,  $q$  and  $\alpha$  (20, 50 and 36 grid points, respectively). We present the best-fitting model, either published or from the grid search, as magenta lines in Figs 5–12. From the  $\chi^2$  confidence interval of the grid-based fits, we estimate a precision of  $\sim 0.2$  dex for the mass ratios. A full binary-lens model will obviously set stronger constraints on the mass ratio in each system. However, for the purpose of our present statistical analysis, the above precision is sufficient.

## 4 PLANET OCCURRENCE FREQUENCY AND MASS-RATIO DISTRIBUTION

The overall abundance of planets near the snowline ( $0.5 < s < 2$ ) can be estimated by integrating over the frequencies of companions up to a certain mass ratio, even without the need to fully constrain the absolute masses of the companions. The population of lenses





**Figure 12.** Inter-calibrated light curves of three anomalous events detected by our detection filter, in addition to those detected both by eye and by the filter, and shown in Figs 5–11. The residuals from the point-lens model are shown in the lower panel of each event, revealing the anomalous region.

in our monitored genII field is dominated by old bulge stars. The mass function of lenses towards the Galactic bulge, as modelled by Dominik (2006), is a narrow distribution, with the most probable lens mass of  $\sim 0.3 M_{\odot}$ . If we define the planet mass limit at  $13 M_J$  (the limiting mass for thermonuclear fusion of deuterium, assuming Solar metallicity, e.g. Perryman 2011), then the corresponding planetary mass-ratio limit is  $q = 4.1 \times 10^{-2}$ . For the present analysis, we will therefore define planets as companions with  $\log q < -1.4$ .

The underlying distribution of mass ratios between companions and primary stars in the population of stars that produce microlensing events can be inferred by correcting the observed distribution of our sample for the detection efficiencies that we derive from the simulations above. The frequency of companions for a given range of mass ratios is simply the number of detected companions within that range,  $N$ , divided by the sum of the detection efficiencies for that range,  $\eta$ , over all events in the sample,

$$f(q) = \frac{N(q)}{\sum \eta(q)}. \quad (5)$$

Fig. 13 shows the observed distribution of mass ratios for our sample in bins of 0.7 dex, and the recovered distribution after accounting for our detection efficiency. The right-hand vertical axis shows the frequencies as a function of mass ratio. We divide the planetary regime into two ranges of mass ratios,  $10^{-2.8} < q < 10^{-1.4}$  and  $10^{-4.9} < q < 10^{-2.8}$ , which correspond to Jupiters and Neptunes, respectively. By choosing these ranges of 1.4 and 2.1 dex, we slightly reduce the Poisson uncertainties, though they still dominate over the detection efficiency uncertainties in each bin. The frequency of snowline Jupiters is

$$f(10^{-2.8} < q < 10^{-1.4}) = 5.0^{+4.0}_{-2.4} \text{ per cent}, \quad (6)$$

and the frequency of Neptunes is

$$f(10^{-4.9} < q < 10^{-2.8}) = 50^{+34}_{-22} \text{ per cent}. \quad (7)$$

Thus, the total frequency of snowline planets of these masses is

$$f(10^{-4.9} < q < 10^{-1.4}) = 55^{+34}_{-22} \text{ per cent}. \quad (8)$$

**Table 1.** Anomalous events. The first 26 events display a clear anomaly and are detected both by eye and by the filter, while the last three, separated by double horizontal lines, are found only by our detection filter.

#	OGLE No.	MOA no.	$q$	Mass from model	Model reference
8	11-0172	11-104	$0.121 \pm 0.005$	$0.02 \pm 0.01 M_{\odot}$	Shin et al. (2012)
10	11-0235	11-107	$\sim 0.77$		
13	11-0265	11-197	$(3.95 \pm 0.06) \times 10^{-3}$	$0.6 \pm 0.2 M_J$	Skowron et al. (2015)
27	11-0481	11-217	$\sim 0.014$		
36	11-0974	11-275	$\sim 0.24$		
46	11-1127	11-322	$0.028 \pm 0.001$		Shvartzvald et al. (2014)
47	11-1132	11-358	$\sim 0.77$		
54	11-9293	11-293	$(5.3 \pm 0.2) \times 10^{-3}$		Yee et al. (2012)
				$4.8 \pm 0.3 M_J$	Batista et al. (2014)
60	12-0442	12-245	$\sim 0.093$		
63	12-0456	12-189	$0.944 \pm 0.004$		Henderson et al. (2014)
96	12-1442	12-532	$0.47 \pm 0.04$		Henderson et al. (2014)
109	13-0341	13-260	$(4.8 \pm 0.3) \times 10^{-5}$	$2.3 \pm 0.2 M_{\oplus}$	Gould et al. (2014)
			$1.21 \pm 0.03$	$0.17 \pm 0.01 M_{\odot}$	
112	13-0506	13-494	$\sim 0.50$		
122	13-0668	13-343	$\sim 0.61$		
137	13-1033	13-472	$\sim 0.4$		
161	13-1721	13-618	$\sim 1.1 \times 10^{-3}$		
167	14-0124	14-307	$(6.9 \pm 0.4) \times 10^{-4}$	$0.5 \pm 0.1 M_J$	Udalski et al. (2015)
168	14-0257	14-148	$\sim 0.24$		
175	14-0572	14-165	$\sim 0.15$		
178	14-0676	14-175	$\sim 1.4 \times 10^{-3}$		
182	14-0727	14-281	$\sim 0.77$		
188	14-0900	14-305	$\sim 0.77$		
191	14-0962	14-285	$\sim 0.98$		
195	14-1109	14-328	$\sim 0.83$		
198	14-1255	14-367	$\sim 0.19$		
222	14-1720	14-507	$\sim 0.093$		
90	12-1211	12-514	$\sim 0.19$		
134	13-0911	13-551	$\sim 2.6 \times 10^{-4}$		
205	14-1406	14-409	$\sim 0.19$		

For the binary regime, we find that the frequency of brown-dwarf companions is

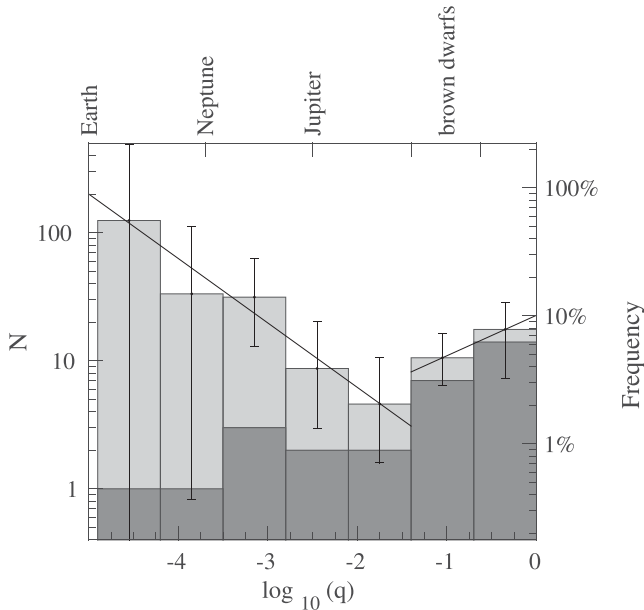
$$f(10^{-1.4} < q < 10^{-0.7}) = 4.7^{+2.6}_{-1.8} \text{ per cent.} \quad (9)$$

and the frequency of stellar companions is

$$f(10^{-0.7} < q < 1) = 7.8^{+4.9}_{-4.6} \text{ per cent.} \quad (10)$$

For a flat intrinsic distribution in  $\log q$ , we expect to find a monotonically rising number of detected companions as a function of  $\log q$ , simply because of the larger cross-section of the more massive planets, and the longer anomaly duration. The recovered mass-ratio distribution in Fig. 13 shows a deficit around  $q \simeq 10^{-2}$ , and the corrected distribution can be fit with a broken – falling and rising – power law. This result echoes previous findings, by radial velocity surveys, that have shown two distinct populations of stellar companions at orbital periods shorter than a few years – planets and stellar binaries – likely produced via two different formation mechanisms (protoplanetary discs, and fragmentation in protostellar clouds, respectively). In those surveys, the two populations are separated by a gap at  $13\text{--}80 M_J$ , the ‘brown-dwarf desert’ (e.g. Grether & Lineweaver 2006). [A recent analysis by Ranc et al. (2015)

compiling current knowledge of brown-dwarf companions from various discovery techniques suggests a more complex brown-dwarf ‘landscape’, with a possible period dependence of brown-dwarf occurrence frequency]. Our detection-efficiency-corrected mass-ratio distribution, assuming our typical primary has a mass of  $0.3 M_{\odot}$ , is suggestive of a similar picture, now possibly seen via microlensing. However, the gap between the two distributions appears to occur at a lower mass,  $\sim 3\text{--}13 M_J$ , corresponding to ‘super-Jupiters’. This difference, if real, could be the result of the fact that microlensing probes the companions of M stars (as opposed to FGK stars by other techniques), that microlensing probes larger separations than other techniques, or both. Due to the small number of total events in our survey, we cannot claim a significant detection of two populations. Furthermore, our mass bins are large, and thus we cannot resolve the region between the putative two populations. Finally, full modelling of the anomalous events, together with additional data (e.g. post-event detection of the lenses) that will help break possible degeneracies in those models, are required in order to give absolute masses rather than only mass ratios. For example, if the hosts of some of the companions in the brown-dwarf bin are in reality more massive than  $0.3 M_{\odot}$ , then their companions will be low-mass M stars. With all of these caveats in mind, the suggestion of a super-Jupiter gap in the M-star companion distribution is nevertheless intriguing. The overall slope we find for the planetary regime ( $10^{-4.9} < q < 10^{-1.4}$ ) is  $d(\log f)/d(\log q) = -0.50 \pm 0.17$ , and for



**Figure 13.** The distribution of estimated mass ratios of our sample, before (dark grey) and after (light grey) correcting for detection efficiency. The right-hand vertical axis shows the corresponding frequency. The top horizontal axis indicates the masses of Earth, Neptune, Jupiter and the brown-dwarf range, for a typical host mass of  $0.3 M_{\odot}$ . The distribution hints at the existence of two populations, stellar binaries and planets, separated by a minimum that is analogous to the brown-dwarf desert found by radial velocity surveys, although here the division appears to be in the super-Jupiter range. Solid lines are power laws  $d(\log f)/d(\log q) = -0.50 \pm 0.17$  (in the planetary regime) and  $d(\log f)/d(\log q) = 0.32 \pm 0.38$  (in the stellar binary regime).

the binary regime ( $10^{-1.4} < q < 1$ ), we find  $d(\log f)/d(\log q) = 0.32 \pm 0.38$ .

## 5 SUMMARY AND COMPARISON TO PREVIOUS WORK

We have presented a statistical analysis of the results, to date, of the first genII microlensing survey. This is the first controlled microlensing experiment for the abundance of snowline planets. Among the 224 events in our sample, 29 are anomalous, revealing the presence of a companion to the lens star. Our results show that about half of the stars have snowline planets with mass ratios in the range  $-4.9 < \log(q) < -1.4$ , with Neptunes being  $\sim 10$  times more common than Jupiters, consistent with previous microlensing findings (e.g. Sumi et al. 2010).

The absolute frequency we find is consistent with previous estimates from first generation microlensing surveys by Gould et al. (2010), who analysed a sample of six detected planets from the follow-up of 13 high-magnification events, and found that the frequency for the range  $-4.5 < \log(q) < -2$  is  $36 \pm 15$  per cent. The frequency we find is somewhat lower than, but consistent with, the estimates by Cassan et al. (2012), who used a control sample of 196 events from the PLANET collaboration with three detected planets (one Jupiter, one Neptune and one super-Earth), and derived a frequency of  $17^{+6}_{-9}$  per cent of Jupiters (in the range of  $0.3\text{--}10 M_J$ ) and  $52^{+22}_{-29}$  per cent Neptunes (in the range of  $10\text{--}30 M_{\oplus}$ ). Our derived planet frequency is also consistent with results from other methods. Radial velocity surveys find that 10.5 per cent of Sun-like

stars host giant planets (in the range of  $0.3\text{--}10 M_J$ ) with periods  $< 5.5$  yr (Cumming et al. 2008). Direct imaging surveys (Brandt et al. 2014), which are currently sensitive only to massive planets ( $> 5 M_J$ ), find that 1.0 per cent  $\text{--}$  3.1 per cent of stars host substellar companions.

For the binary regime, we have estimated from our survey that  $\sim 8$  per cent of lens stars have stellar companions with separations in the range  $0.5 < s < 2$ . This is consistent with direct imaging estimates from the M-dwarfs in Multiples survey, of  $8 \pm 2$  per cent for companions to M stars with separations of  $3\text{--}10$  AU (Ward-Duong et al. 2015), and with radial-velocity surveys estimates of  $\sim 10$  per cent for companions to FGK stars with periods  $< 5$  yr (Grether & Lineweaver 2006; Raghavan et al. 2010). Although we cannot clearly identify the brown-dwarf regime, we can estimate the frequency of brown-dwarf companions at  $\sim 5$  per cent. This is significantly higher than the estimates based on radial-velocity surveys of FGK stars, of  $< 1$  per cent, by Grether & Lineweaver (2006) for companions with periods  $< 5$  yr. Similarly, the frequency of brown-dwarf companions to white dwarfs found by Steele et al. (2011) is also low,  $f_{WD+BD} = 0.5 \text{ per cent} \pm 0.3 \text{ per cent}$ . Typical white dwarfs, with masses of  $0.6 M_{\odot}$ , are descended from  $\sim 2 M_{\odot}$  stars. It is presently unclear whether brown-dwarf companions become more common with increasing separation. For example, Metchev & Hillenbrand (2009) estimate a frequency of  $3.2^{+3.1}_{-2.7}$  per cent of brown-dwarf companions in  $28\text{--}1590$  AU orbits around Sun-like stars. This fraction is consistent with all the various brown-dwarf companion estimates above.

The companion mass-ratio distribution in our sample shows a decline from companions with similar masses ( $q \simeq 1$ ) towards companions with  $q \simeq 10^{-2}$ . After accounting for the detection efficiency for different mass ratios, the distribution is suggestive of a rise in the planetary regime from massive planets towards less massive planets. As discussed above, the minimum in the distribution may be the microlensing manifestation of the bimodal companion distribution found by other techniques, but with the minimum shifted from the brown-dwarf mass range to the super-Jupiter range.

The slope we find for the planetary regime is consistent with what has been found for planetary companions to FGK stars from radial velocity surveys. For example, Cumming et al. (2008) find a logarithmic slope of  $d(\log f)/d(\log M) = -0.31 \pm 0.2$  for companions in the mass range of  $0.3\text{--}10 M_J$  and with orbital periods of  $2\text{--}2000$  d. It is also consistent with direct-imaging exoplanet survey results, e.g. Brandt et al. (2014), who find a slope of  $d(\log f)/d(\log M) = -0.65 \pm 0.60$  for the distribution of substellar companions with masses of  $5\text{--}70 M_J$  between  $10$  and  $100$  AU. For the binary regime, our slope is shallower but, again, consistent with the  $d(\log f)/d(\log M) = 0.68 \pm 0.21$  found by radial velocity surveys (Grether & Lineweaver 2006) for companions with masses of  $0.1 < (M/M_{\odot}) < 1$ , and with the  $d(\log f)/d(\log M) \sim 0.5$  shown by Mazeh et al. (2003) for the companion distribution in that mass range.

Full binary-lens modelling of the anomalous events in our sample is required in order to confirm these estimates. For some of the events, the light curves show the signatures of high-order effects, and thus the mass of the host star and its companion can be measured. For the rest of the sample, high-resolution imaging in the coming years could refine estimates of the masses, by isolating the light from the lens. Together with increasing numbers of planetary detections from our ongoing genII survey, this will give a progressively clearer picture of the planetary mass distribution near the snowlines of low-mass stars in the Galaxy.

## ACKNOWLEDGEMENTS

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## APPENDIX A

Microlens event sample: best-fitting point-lens model parameters. Significant detections of finite source effects and parallax are marked in bold.

At the end of the table, separated by double horizontal lines, we list the six anomalous events which passed the limiting magnitude criteria only because of the anomaly, and are therefore excluded from the sample (see Section 2.2 for details).

**Table A1.** Event summary.

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E,E}$	$\pi_{E,N}$	$I_{bl}$	$f_{bl}$
1	11-0022	11-025	0.867 (0.002)	5690.508 (0.041)	56.584 (0.150)	0.000 (38.685)	<b>-0.379</b> <b>(0.003)</b>	<b>0.021</b> <b>(0.007)</b>	15.614 (0.000)	0.000 (0.000)
2	11-0037	11-039	0.173 (0.001)	5807.481 (0.025)	128.572 (0.565)	<b>2.284</b> <b>(14.211)</b>	<b>0.263</b> <b>(0.001)</b>	<b>-0.243</b> <b>(0.002)</b>	16.155 (0.000)	0.245 (0.001)
3	11-0081	11-064	0.207 (0.003)	5658.523 (0.128)	106.779 (1.568)	<b>223.877</b> <b>(3.383)</b>	<b>0.437</b> <b>(0.014)</b>	<b>0.457</b> <b>(0.011)</b>	16.460 (0.000)	0.952 (0.000)
4	11-0082	11-103	0.295 (0.009)	5702.214 (0.037)	63.944 (1.186)	169.666 (49.256)	-0.205 (0.030)	-0.464 (0.058)	18.191 (0.001)	0.057 (0.002)
5	11-0120	11-114	0.846 (0.031)	5670.592 (0.109)	26.837 (0.698)	0.000 (71.112)	0.015 (0.040)	1.496 (0.044)	16.964 (0.001)	0.003 (0.004)
6	11-0138	11-082	0.042 (0.001)	5666.781 (0.008)	52.366 (0.920)	6.039 (5.511)	0.206 (0.027)	-0.203 (0.143)	19.672 (0.002)	0.251 (0.002)
7	11-0168	11-091	0.351 (0.004)	5690.133 (0.014)	24.810 (0.160)	0.000 (45.810)	0.158 (0.025)	1.498 (0.104)	16.550 (0.000)	0.215 (0.001)
8	11-0172	11-104	-0.050 (0.001)	5670.367 (0.010)	49.263 (0.918)	51.767 (1.179)	-0.046 (0.030)	-1.493 (0.068)	19.497 (0.003)	0.113 (0.003)
9	11-0173	11-133	-0.696 (0.052)	5689.328 (0.059)	30.654 (1.142)	0.000 (147.475)	<b>-0.160</b> <b>(0.029)</b>	<b>-1.315</b> <b>(0.201)</b>	15.828 (0.000)	0.438 (0.001)
10	11-0235	11-107	-0.669 (0.050)	5673.096 (0.172)	14.184 (0.664)	0.000 (143.850)	0.220 (0.418)	1.208 (0.811)	18.155 (0.002)	0.000 (0.016)
11	11-0240	11-109	0.035 (0.018)	5659.760 (0.023)	9.381 (1.604)	44.823 (24.768)	1.424 (0.802)	-1.084 (0.793)	19.000 (0.004)	0.939 (0.001)
12	11-0243	11-151	0.700 (0.011)	5714.091 (0.057)	26.304 (0.243)	0.000 (98.704)	0.461 (0.068)	1.499 (0.068)	16.707 (0.000)	0.000 (0.000)
13	11-0265	11-197	-0.136 (0.000)	5760.259 (0.009)	51.582 (0.152)	0.000 (10.814)	0.061 (0.017)	0.628 (0.110)	17.287 (0.000)	0.000 (0.000)
14	11-0266	11-115	0.664 (0.007)	5694.409 (0.012)	14.482 (0.117)	0.000 (71.339)	0.489 (0.088)	0.858 (0.373)	15.559 (0.000)	0.000 (0.001)
15	11-0286	11-124	0.109 (0.032)	5680.101 (0.012)	12.338 (0.496)	76.999 (28.270)	1.494 (0.340)	-0.978 (0.788)	18.636 (0.003)	0.620 (0.002)
16	11-0293	11-190	-0.476 (0.026)	5714.295 (0.052)	32.685 (1.163)	414.080 (119.996)	-0.037 (0.151)	1.483 (0.941)	18.501 (0.002)	0.033 (0.004)
17	11-0297	11-166	0.015 (0.002)	5700.670 (0.021)	136.964 (14.527)	19.059 (4.812)	-0.605 (0.057)	-1.446 (0.096)	18.794 (0.004)	0.907 (0.001)
18	11-0318	11-132	0.503 (0.106)	5682.097 (0.002)	1.055 (0.023)	356.525 (131.754)	1.312 (0.804)	-1.171 (0.817)	16.002 (0.001)	0.006 (0.003)
19	11-0329	11-147	-0.005 (0.000)	5691.899 (0.000)	11.928 (0.231)	5.400 (0.297)	-0.663 (0.541)	-0.720 (0.797)	19.194 (0.004)	0.255 (0.003)
20	11-0348	11-189	0.443 (0.038)	5713.657 (0.041)	19.410 (0.423)	85.249 (111.233)	0.674 (0.446)	1.484 (0.828)	18.349 (0.001)	0.073 (0.004)
21	11-0364	11-154	0.019 (0.006)	5703.622 (0.009)	74.999 (6.046)	17.449 (5.385)	<b>-0.798</b> <b>(0.068)</b>	<b>-0.242</b> <b>(0.061)</b>	16.101 (0.000)	0.986 (0.000)
22	11-0421	11-160	0.050 (0.003)	5705.967 (0.004)	7.771 (0.262)	0.583 (12.883)	-1.490 (0.802)	1.465 (0.813)	16.247 (0.000)	0.940 (0.000)
23	11-0422	11-171	0.207 (0.013)	5709.443 (0.026)	17.937 (0.848)	218.483 (14.174)	1.490 (0.099)	-1.496 (0.336)	18.477 (0.002)	0.537 (0.002)
24	11-0424	11-204	0.314 (0.053)	5732.063 (0.011)	16.901 (0.125)	193.531 (73.532)	-1.480 (0.270)	0.330 (0.300)	16.299 (0.001)	0.124 (0.002)
25	11-0462	11-191	0.004 (0.000)	5763.322 (0.000)	179.829 (0.618)	2.950 (0.143)	<b>-0.068</b> <b>(0.004)</b>	<b>0.098</b> <b>(0.012)</b>	16.406 (0.000)	0.935 (0.000)
26	11-0465	11-211	0.562 (0.005)	5781.460 (0.016)	35.139 (0.158)	109.840 (68.034)	<b>-0.127</b> <b>(0.016)</b>	<b>0.582</b> <b>(0.054)</b>	14.944 (0.000)	0.000 (0.001)
27	11-0481	11-217	-0.018 (0.016)	5725.269 (0.003)	5.607 (0.389)	2.650 (17.413)	0.873 (0.808)	0.515 (0.792)	19.111 (0.003)	0.601 (0.002)
28	11-0507	11-287	0.252 (0.021)	5750.747 (0.055)	35.440 (2.027)	0.000 (57.304)	<b>-0.974</b> <b>(0.362)</b>	<b>1.482</b> <b>(0.670)</b>	17.153 (0.001)	0.867 (0.001)
29	11-0510	11-247	0.664 (0.013)	5741.451 (0.010)	8.633 (0.101)	270.825 (104.230)	-1.433 (0.541)	-1.496 (0.568)	16.884 (0.000)	0.000 (0.000)
30	11-0519	11-273	0.257 (0.027)	5741.040 (0.019)	10.927 (0.560)	0.000 (76.924)	1.479 (0.718)	1.499 (0.799)	15.615 (0.000)	0.943 (0.000)
31	11-0521	11-270	1.025 (0.033)	5746.899 (0.016)	7.426 (0.173)	38.738 (52.255)	-1.498 (0.106)	-1.491 (0.178)	14.231 (0.000)	0.001 (0.002)
32	11-0535	11-354	-0.083 (0.008)	5768.233 (0.040)	73.757 (5.729)	82.167 (23.369)	-0.038 (0.045)	0.084 (0.250)	18.321 (0.001)	0.914 (0.000)



Table A1 – continued

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E,E}$	$\pi_{E,N}$	$I_{bl}$	$f_{bl}$
33	11-0545	11-290	0.238 (0.064)	5751.104 (0.045)	17.559 (0.360)	0.000 (60.034)	-0.205 (0.564)	-1.447 (0.795)	19.099 (0.004)	0.001 (0.008)
34	11-0944	11-297	-0.046 (0.000)	5767.394 (0.001)	12.273 (0.013)	1.624 (5.479)	0.288 (0.144)	-1.499 (0.654)	16.983 (0.000)	0.000 (0.000)
35	11-0966	11-302	0.077 (0.020)	5758.121 (0.016)	16.523 (0.347)	17.351 (21.583)	0.128 (0.419)	-1.427 (0.836)	19.114 (0.002)	0.137 (0.004)
36	11-0974	11-275	0.004 (18.946)	5743.417 (0.092)	267.266 (122.304)	4.783 (16117.873)	-0.981 (0.802)	-1.404 (0.795)	19.485 (0.455)	0.995 (0.006)
37	11-0990	11-300	-0.012 (0.003)	5758.680 (0.004)	6.788 (0.108)	16.325 (4.291)	-0.975 (0.736)	-0.991 (0.797)	18.480 (0.002)	0.044 (0.005)
38	11-0991	11-312	0.810 (0.033)	5801.829 (0.049)	25.457 (0.595)	192.278 (116.572)	0.320 (0.026)	-1.500 (0.083)	16.018 (0.000)	0.174 (0.001)
39	11-0999	11-306	0.109 (0.021)	5760.856 (0.006)	5.487 (0.255)	44.185 (33.615)	1.419 (0.802)	1.178 (0.800)	19.616 (0.003)	0.197 (0.005)
40	11-1003	11-298	0.071 (0.039)	5757.933 (0.075)	5.957 (0.384)	161.260 (47.824)	-1.142 (0.811)	-0.224 (0.818)	17.404 (0.001)	0.759 (0.005)
41	11-1007	11-370	-0.273 (0.008)	5789.990 (0.053)	34.525 (0.836)	121.625 (56.154)	-0.045 (0.053)	-1.435 (0.269)	18.798 (0.002)	0.003 (0.005)
42	11-1009	11-335	0.022 (0.006)	5774.069 (0.058)	180.284 (35.620)	6.455 (6.526)	0.029 (0.103)	-0.025 (0.641)	16.259 (0.002)	0.994 (0.000)
43	11-1035	11-337	-0.052 (0.005)	5776.564 (0.010)	32.138 (1.419)	34.412 (17.005)	-0.343 (0.159)	0.032 (0.521)	19.026 (0.002)	0.796 (0.001)
44	11-1036	11-372	-0.198 (0.018)	5791.009 (0.048)	34.887 (1.452)	2.241 (69.196)	0.027 (0.066)	0.025 (0.671)	18.932 (0.003)	0.230 (0.005)
45	11-1072	11-346	0.006 (0.000)	5775.398 (0.001)	15.875 (0.293)	5.552 (0.458)	-1.496 (0.140)	-0.263 (0.715)	18.143 (0.002)	0.536 (0.001)
46	11-1127	11-322	-0.005 (0.010)	5775.101 (0.010)	23.687 (2.041)	90.657 (7.448)	1.494 (0.124)	1.495 (0.737)	16.699 (0.001)	0.945 (0.001)
47	11-1132	11-358	0.018 (0.003)	5753.381 (0.073)	298.017 (73.846)	17.452 (5.238)	0.721 (0.154)	-0.529 (0.224)	18.800 (0.003)	0.970 (0.000)
48	11-1162	11-379	0.693 (0.035)	5799.155 (0.037)	10.703 (0.349)	1.114 (128.940)	-0.380 (0.286)	-1.497 (0.808)	17.437 (0.001)	0.089 (0.004)
49	11-1195	11-389	-0.481 (0.092)	5804.265 (0.110)	23.892 (2.200)	459.210 (132.287)	0.055 (0.141)	1.376 (0.832)	18.218 (0.009)	0.150 (0.015)
50	11-1200	11-381	-0.081 (0.006)	5799.188 (0.006)	13.837 (0.497)	83.545 (8.546)	-0.061 (0.364)	0.207 (0.780)	17.687 (0.001)	0.740 (0.002)
51	11-1254	11-388	-0.006 (0.002)	5801.312 (0.003)	29.373 (4.858)	<b>10.393</b> ( <b>1.660</b> )	0.737 (0.482)	1.372 (0.803)	15.934 (0.000)	0.996 (0.000)
52	11-9112	11-112	-0.310 (0.012)	5686.106 (0.050)	20.708 (0.469)	0.000 (63.919)	<b>0.427</b> ( <b>0.097</b> )	<b>1.494</b> ( <b>0.749</b> )	16.368 (0.002)	0.000 (0.000)
53	11-9291	11-291	0.003 (0.001)	5747.950 (0.001)	58.011 (2.370)	1.766 (0.749)	-1.494 (0.100)	-1.490 (0.250)	17.420 (0.001)	0.988 (0.000)
54	11-9293	11-293	0.001 (0.000)	5747.502 (0.002)	79.317 (6.984)	1.498 (0.174)	-1.457 (0.389)	-1.462 (0.689)	19.935 (0.016)	0.973 (0.000)
55	11-9313	11-313	-0.189 (0.017)	5766.639 (0.008)	5.284 (0.357)	211.722 (18.439)	1.497 (0.735)	0.798 (0.789)	18.121 (0.007)	0.640 (0.013)
56	11-9367	11-367	0.052 (0.003)	5787.101 (0.002)	9.833 (0.202)	3.286 (16.151)	1.452 (0.399)	1.316 (0.804)	18.770 (0.002)	0.314 (0.002)
57	11-9393	11-393	0.001 (0.002)	5803.567 (0.002)	20.698 (1.499)	5.645 (1.537)	0.414 (0.457)	0.341 (0.767)	19.005 (0.004)	0.878 (0.001)
58	12-0305	12-182	-0.685 (0.025)	6046.970 (0.046)	20.933 (0.611)	402.896 (123.591)	-0.356 (0.057)	0.217 (0.575)	16.333 (0.001)	0.207 (0.002)
59	12-0325	12-166	0.972 (0.041)	6039.774 (0.092)	15.785 (0.808)	2.207 (113.530)	0.609 (0.175)	1.485 (0.886)	16.651 (0.001)	0.000 (0.000)
60	12-0442	12-245	0.174 (0.005)	6078.680 (0.017)	31.894 (0.541)	105.718 (32.295)	-0.344 (0.051)	1.499 (0.026)	18.214 (0.002)	0.164 (0.003)
61	12-0443	12-211	0.058 (0.002)	6046.093 (0.011)	24.627 (0.705)	41.639 (14.238)	-0.342 (0.164)	-0.109 (0.667)	19.515 (0.005)	0.215 (0.005)
62	12-0449	12-216	-0.725 (0.105)	6058.735 (0.089)	35.066 (2.285)	706.358 (222.866)	-0.336 (0.074)	1.497 (0.965)	16.646 (0.000)	0.812 (0.001)
63	12-0456	12-189	-0.169 (0.008)	6047.058 (0.013)	7.031 (0.019)	9.316 (35.654)	-1.494 (0.070)	1.422 (0.573)	15.179 (0.000)	0.000 (0.000)
64	12-0462	12-271	0.055 (0.000)	6064.302 (0.006)	37.484 (0.186)	0.000 (6.292)	-0.014 (0.049)	1.499 (0.083)	19.046 (0.001)	0.000 (0.000)

Table A1 – *continued*

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E, E}$	$\pi_{E, N}$	$I_{bl}$	$f_{bl}$
65	12-0591	12-430	0.175 (0.003)	6170.359 (0.012)	63.396 (0.491)	78.360 (24.395)	−0.013 (0.006)	−0.245 (0.080)	17.075 (0.001)	0.247 (0.001)
66	12-0615	12-277	0.507 (0.021)	6060.504 (0.007)	4.098 (0.090)	1.809 (77.143)	−1.492 (0.412)	1.438 (0.772)	14.411 (0.000)	0.753 (0.001)
67	12-0694	12-308	0.042 (0.014)	6070.255 (0.007)	12.163 (1.535)	11.372 (11.920)	1.372 (0.761)	−0.673 (0.773)	19.593 (0.006)	0.744 (0.003)
68	12-0722	12-397	−0.173 (0.011)	6122.064 (0.066)	58.003 (2.921)	9.264 (46.579)	−0.067 (0.115)	−1.303 (0.997)	18.345 (0.005)	0.501 (0.005)
69	12-0724	12-323	0.009 (0.001)	6071.042 (0.002)	15.176 (2.544)	<b>11.708</b> (1.885)	−1.385 (0.756)	1.047 (0.749)	19.821 (0.005)	0.880 (0.001)
70	12-0726	12-351	−0.204 (0.075)	6072.873 (0.030)	11.110 (0.644)	1.535 (71.483)	1.489 (0.478)	−1.473 (0.806)	18.183 (0.003)	0.503 (0.004)
71	12-0784	12-337	−0.029 (0.005)	6085.181 (0.018)	88.027 (24.624)	27.140 (6.748)	−0.061 (0.084)	−0.027 (0.252)	16.176 (0.000)	0.988 (0.000)
72	12-0825	12-451	0.352 (0.010)	6171.507 (0.036)	45.769 (2.222)	8.244 (48.109)	0.073 (0.021)	0.314 (0.204)	17.399 (0.001)	0.002 (0.002)
73	12-0867	12-352	0.202 (0.001)	6110.261 (0.002)	13.949 (0.035)	0.000 (18.564)	0.128 (0.242)	−1.358 (0.420)	14.260 (0.000)	0.045 (0.001)
74	12-1414	12-205	−0.043 (0.003)	6040.895 (0.011)	36.153 (1.728)	0.000 (11.665)	−0.195 (0.127)	−0.186 (0.667)	20.257 (0.007)	0.280 (0.006)
75	12-1430	12-278	−0.002 (0.001)	6062.104 (0.001)	18.807 (0.268)	4.130 (1.150)	0.525 (0.345)	0.941 (0.835)	20.217 (0.006)	0.001 (0.006)
76	12-0941	12-446	0.245 (0.011)	6126.169 (0.031)	36.883 (1.256)	0.000 (50.847)	−0.537 (0.238)	1.499 (0.726)	19.133 (0.003)	0.113 (0.004)
77	12-0974	12-424	−0.069 (0.019)	6110.120 (0.079)	86.425 (11.899)	91.769 (9.884)	−1.477 (0.143)	−0.031 (0.209)	18.802 (0.005)	0.939 (0.001)
78	12-0998	12-447	−0.518 (0.036)	6136.857 (0.033)	24.356 (0.513)	0.000 (179.737)	0.125 (0.069)	−1.439 (0.881)	16.957 (0.000)	0.482 (0.001)
79	12-1013	12-449	0.643 (0.108)	6126.889 (0.078)	21.525 (2.017)	630.572 (197.902)	−0.858 (0.286)	0.248 (0.627)	18.599 (0.002)	0.209 (0.006)
80	12-1049	12-438	−0.338 (0.098)	6117.274 (0.017)	7.192 (0.309)	152.796 (95.418)	−1.486 (0.732)	−1.069 (0.802)	18.302 (0.002)	0.371 (0.003)
81	12-1051	12-477	0.200 (0.026)	6132.237 (0.014)	21.209 (0.273)	0.000 (37.963)	0.025 (0.115)	−0.565 (0.741)	18.494 (0.002)	0.017 (0.004)
82	12-1418	12-435	0.018 (0.001)	6119.272 (0.003)	55.686 (3.094)	0.000 (3.930)	0.280 (0.209)	−0.016 (0.518)	20.142 (0.006)	0.539 (0.003)
83	12-1067	12-444	−0.011 (0.001)	6116.566 (0.000)	7.158 (0.470)	0.396 (2.402)	−1.486 (0.797)	−1.430 (0.816)	18.966 (0.003)	0.848 (0.001)
84	12-1069	12-452	−0.220 (0.006)	6129.142 (0.020)	22.575 (0.199)	0.913 (51.679)	0.281 (0.173)	−1.495 (0.765)	19.044 (0.002)	0.000 (0.000)
85	12-1066	12-476	−0.132 (0.002)	6148.598 (0.004)	17.562 (0.096)	80.624 (19.811)	−0.086 (0.050)	0.405 (0.548)	16.670 (0.001)	0.048 (0.001)
86	12-1074	12-519	−0.259 (0.013)	6154.066 (0.050)	40.088 (1.301)	183.433 (60.351)	0.215 (0.050)	0.052 (0.579)	19.280 (0.003)	0.001 (0.005)
87	12-1183	12-478	−0.107 (0.015)	6130.447 (0.028)	7.131 (0.611)	128.561 (38.220)	0.926 (0.807)	1.082 (0.818)	19.427 (0.008)	0.593 (0.008)
88	12-1193	12-498	0.106 (0.008)	6136.399 (0.026)	14.384 (0.884)	118.142 (9.552)	1.295 (0.527)	1.381 (0.780)	18.296 (0.002)	0.801 (0.002)
89	12-1210	12-511	−0.368 (0.010)	6143.702 (0.006)	1.881 (0.015)	1.844 (82.426)	1.499 (0.796)	1.498 (0.812)	16.507 (0.000)	0.000 (0.000)
90	12-1211	12-514	−0.614 (0.010)	6151.519 (0.011)	6.542 (0.086)	621.011 (10.763)	1.489 (0.081)	1.070 (0.799)	15.699 (0.000)	0.332 (0.003)
91	12-1213	12-512	0.290 (0.045)	6144.519 (0.057)	18.487 (0.570)	327.066 (111.028)	−1.476 (0.175)	−1.420 (0.744)	19.538 (0.005)	0.000 (0.000)
92	12-1250	12-515	0.060 (0.013)	6145.218 (0.022)	6.409 (0.897)	96.900 (29.045)	1.285 (0.792)	0.452 (0.793)	18.184 (0.002)	0.862 (0.004)
93	12-1242	12-552	−0.059 (0.003)	6160.797 (0.009)	25.921 (0.863)	1.650 (13.699)	−0.380 (0.130)	0.219 (0.724)	19.306 (0.002)	0.581 (0.002)
94	12-1244	12-555	0.214 (0.066)	6154.529 (0.093)	38.326 (6.447)	231.692 (86.241)	0.227 (0.165)	−0.978 (0.871)	19.626 (0.006)	0.534 (0.006)
95	12-1245	12-575	−0.229 (0.028)	6173.239 (0.137)	41.186 (4.094)	0.000 (63.580)	0.056 (0.101)	0.125 (0.660)	19.264 (0.005)	0.320 (0.007)
96	12-1442	12-532	0.025 (0.035)	6151.813 (0.017)	11.310 (1.859)	6.815 (27.043)	−0.388 (0.770)	−0.053 (0.800)	20.194 (0.018)	0.242 (0.018)

Table A1 – continued

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E,E}$	$\pi_{E,N}$	$I_{bl}$	$f_{bl}$
97	12-1268	12-574	0.204 (0.004)	6171.071 (0.021)	21.705 (0.294)	0.311 (43.168)	0.259 (0.086)	-1.498 (0.805)	18.668 (0.002)	0.000 (0.003)
98	12-1269	12-560	0.217 (0.021)	6160.201 (0.026)	12.653 (0.757)	202.708 (48.311)	-1.066 (0.468)	0.733 (0.816)	18.311 (0.002)	0.733 (0.002)
99	12-1292	12-570	0.214 (0.054)	6165.816 (0.003)	5.658 (0.067)	89.066 (44.433)	-0.262 (0.500)	0.688 (0.795)	16.986 (0.000)	0.082 (0.002)
100	12-1581	12-577	-0.149 (0.007)	6169.822 (0.043)	35.435 (1.747)	118.073 (36.750)	-0.091 (0.087)	-1.424 (0.986)	20.100 (0.005)	0.027 (0.007)
101	12-1311	12-594	0.588 (0.127)	6170.085 (0.172)	20.773 (2.486)	0.000 (216.333)	1.056 (0.260)	0.782 (0.811)	17.841 (0.001)	0.646 (0.004)
102	12-1364	12-580	-0.015 (0.005)	6164.630 (0.003)	16.616 (0.987)	0.000 (3.781)	1.399 (0.314)	1.491 (0.815)	18.465 (0.003)	0.739 (0.001)
103	12-9582	12-582	-0.015 (0.012)	6162.944 (0.002)	1.920 (4.803)	30.483 (15.692)	-0.189 (0.794)	-0.202 (0.798)	18.060 (0.011)	0.929 (0.027)
104	12-9591	12-591	0.006 (0.003)	6166.857 (0.009)	18.228 (5.984)	6.773 (3.041)	0.702 (0.785)	-0.735 (0.788)	18.644 (0.012)	0.991 (0.001)
105	12-1350	12-603	-0.134 (0.011)	6175.036 (0.018)	10.578 (0.474)	136.409 (37.325)	-1.241 (0.534)	0.313 (0.820)	18.846 (0.006)	0.084 (0.009)
106	13-0036	13-399	0.617 (0.011)	6406.163 (0.036)	47.908 (0.393)	0.000 (101.668)	0.040 (0.010)	0.141 (0.041)	15.971 (0.000)	0.038 (0.001)
107	13-0145	13-110	0.467 (0.007)	6405.939 (0.019)	29.063 (0.226)	0.065 (54.737)	-0.009 (0.024)	0.562 (0.233)	16.310 (0.000)	0.000 (0.002)
108	13-0183	13-332	-0.652 (0.023)	6495.203 (0.135)	103.843 (2.225)	15.347 (91.360)	<b>-0.076</b> <b>(0.006)</b>	<b>0.035</b> <b>(0.009)</b>	16.888 (0.000)	0.403 (0.001)
109	13-0341	13-260	0.000 (0.005)	6406.313 (0.005)	28.880 (0.335)	13.446 (5.332)	1.063 (0.063)	-0.909 (0.493)	18.368 (0.002)	0.066 (0.002)
110	13-0486	13-294	-0.154 (0.004)	6420.492 (0.015)	28.978 (0.389)	0.456 (33.668)	-0.171 (0.055)	-0.062 (0.367)	18.247 (0.001)	0.480 (0.001)
111	13-0488	13-355	-0.011 (0.004)	6439.171 (0.005)	93.899 (4.458)	0.720 (2.988)	-0.173 (0.121)	-0.448 (0.372)	19.738 (0.007)	0.517 (0.004)
112	13-0506	13-494	-0.089 (0.036)	6501.345 (0.246)	154.414 (13.158)	177.820 (39.844)	-0.057 (0.015)	-0.011 (0.099)	17.570 (0.001)	0.898 (0.001)
113	13-0513	13-272	0.136 (0.010)	6407.452 (0.038)	28.801 (1.433)	0.000 (37.790)	0.502 (0.145)	1.481 (0.854)	19.766 (0.005)	0.202 (0.006)
114	13-0601	13-292	-0.060 (0.015)	6414.867 (0.016)	17.778 (1.414)	80.207 (22.563)	-1.158 (0.390)	0.987 (0.828)	18.756 (0.003)	0.810 (0.002)
115	13-0607	13-376	-0.956 (0.024)	6455.457 (0.044)	22.750 (0.434)	0.000 (141.951)	-0.189 (0.171)	1.492 (0.885)	16.383 (0.000)	0.001 (0.002)
116	13-0608	13-308	-0.438 (0.088)	6426.720 (0.023)	14.260 (0.310)	378.693 (102.202)	0.559 (0.224)	0.521 (0.736)	17.388 (0.001)	0.023 (0.003)
117	13-0610	13-401	0.253 (0.002)	6455.122 (0.004)	22.810 (0.077)	0.000 (28.019)	1.084 (0.087)	1.499 (0.106)	16.059 (0.000)	0.071 (0.001)
118	13-0615	13-389	-0.349 (0.036)	6448.248 (0.117)	46.213 (2.467)	0.484 (113.663)	0.248 (0.129)	1.500 (0.828)	18.110 (0.002)	0.484 (0.004)
119	13-0616	13-331	-0.355 (0.084)	6424.294 (0.019)	10.334 (0.377)	0.000 (97.568)	-0.841 (0.430)	1.490 (0.775)	17.882 (0.001)	0.381 (0.002)
120	13-0619	13-351	-0.055 (0.003)	6427.437 (0.007)	21.427 (0.708)	0.000 (14.559)	0.571 (0.330)	1.497 (0.839)	20.215 (0.006)	0.169 (0.006)
121	13-0621	13-344	-0.101 (0.010)	6425.046 (0.019)	30.293 (1.718)	85.822 (32.894)	-0.290 (0.188)	1.496 (0.848)	19.606 (0.003)	0.610 (0.002)
122	13-0668	13-343	0.246 (0.068)	6422.307 (0.058)	22.570 (2.768)	266.126 (82.594)	0.806 (0.348)	0.020 (0.771)	19.287 (0.003)	0.693 (0.004)
123	13-0674	13-346	-0.002 (0.000)	6426.021 (0.000)	57.315 (1.205)	0.000 (0.251)	-0.194 (0.194)	-1.473 (0.880)	20.105 (0.006)	0.738 (0.002)
124	13-0682	13-388	0.727 (0.013)	6453.990 (0.045)	21.275 (0.231)	78.987 (103.167)	1.159 (0.281)	1.488 (0.256)	16.880 (0.001)	0.002 (0.003)
125	13-0689	13-379	0.237 (0.031)	6434.013 (0.099)	21.510 (1.625)	8.974 (76.931)	-0.162 (0.477)	1.494 (0.841)	17.949 (0.002)	0.814 (0.002)
126	13-0703	13-349	-0.151 (0.044)	6429.186 (0.019)	25.032 (0.740)	81.227 (28.229)	0.129 (0.170)	-0.119 (0.650)	19.014 (0.003)	0.146 (0.004)
127	13-0731	13-377	-0.292 (0.073)	6437.552 (0.058)	18.103 (1.011)	0.000 (90.446)	-0.247 (0.463)	0.149 (0.774)	17.346 (0.001)	0.804 (0.001)
128	13-0770	13-462	0.238 (0.008)	6471.819 (0.023)	26.446 (0.508)	98.029 (49.735)	<b>-1.499</b> <b>(0.026)</b>	<b>-0.862</b> <b>(0.178)</b>	17.449 (0.001)	0.342 (0.002)

Table A1 – *continued*

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E, E}$	$\pi_{E, N}$	$I_{bl}$	$f_{bl}$
129	13-0801	13-386	0.047 (0.008)	6436.617 (0.007)	11.463 (1.440)	56.062 (9.493)	−1.473 (0.769)	1.028 (0.780)	18.103 (0.003)	0.949 (0.001)
130	13-0835	13-400	−0.056 (0.001)	6449.971 (0.001)	9.743 (0.092)	0.000 (9.652)	−1.489 (0.692)	−0.011 (0.781)	18.184 (0.002)	0.204 (0.002)
131	13-0850	13-470	0.959 (0.070)	6479.565 (0.072)	22.525 (1.258)	0.486 (152.823)	0.078 (0.280)	−0.607 (0.352)	17.165 (0.001)	0.007 (0.003)
132	13-0861	13-502	0.363 (0.017)	6516.004 (0.101)	65.655 (1.497)	0.000 (87.172)	0.037 (0.020)	−0.502 (0.079)	17.402 (0.001)	0.324 (0.002)
133	13-0891	13-433	−0.609 (0.003)	6476.835 (0.011)	18.930 (0.092)	0.000 (54.934)	0.322 (0.105)	−0.737 (0.473)	15.512 (0.000)	0.000 (0.000)
134	13-0911	13-551	−0.003 (0.000)	6537.299 (0.000)	126.734 (1.934)	0.533 (0.720)	−0.131 (0.004)	0.218 (0.077)	19.362 (0.003)	0.444 (0.002)
135	13-0925	13-493	−0.358 (0.015)	6501.455 (0.055)	43.130 (0.977)	171.945 (81.057)	−0.194 (0.027)	1.494 (0.081)	18.051 (0.001)	0.169 (0.003)
136	13-0968	13-452	−0.168 (0.004)	6460.311 (0.002)	7.046 (0.056)	21.913 (32.723)	−1.489 (0.783)	1.279 (0.750)	17.675 (0.001)	0.049 (0.001)
137	13-1033	13-472	−0.139 (0.098)	6468.808 (0.050)	21.183 (3.002)	196.985 (82.722)	−0.371 (0.777)	1.377 (0.760)	19.571 (0.005)	0.550 (0.007)
138	13-1037	13-483	0.048 (0.015)	6473.098 (0.019)	36.512 (3.706)	1.531 (16.175)	−1.500 (0.511)	−0.754 (0.697)	20.185 (0.013)	0.437 (0.009)
139	13-1078	13-460	0.000 (0.000)	6462.061 (0.000)	27.261 (1.700)	<b>0.990</b> ( <b>0.072</b> )	−1.210 (0.798)	1.338 (0.806)	20.259 (0.015)	0.929 (0.002)
140	13-1080	13-456	0.004 (0.002)	6462.038 (0.001)	22.224 (4.915)	0.697 (1.221)	1.380 (0.808)	−1.492 (0.818)	20.388 (0.007)	0.904 (0.001)
141	13-9455	13-455	0.058 (0.030)	6456.795 (0.023)	7.492 (1.567)	110.370 (32.662)	0.588 (0.811)	−1.359 (0.812)	18.519 (0.012)	0.684 (0.019)
142	13-1114	13-485	0.008 (0.001)	6474.942 (0.003)	68.566 (5.160)	0.000 (2.113)	−0.707 (0.147)	−0.383 (0.300)	20.374 (0.018)	0.292 (0.012)
143	13-1124	13-500	0.269 (0.047)	6489.987 (0.045)	22.335 (1.035)	0.727 (96.262)	0.671 (0.263)	−1.478 (0.806)	18.292 (0.001)	0.593 (0.002)
144	13-1145	13-523	−0.053 (0.011)	6503.507 (0.009)	36.562 (1.204)	8.814 (13.386)	−0.173 (0.131)	1.443 (0.886)	18.536 (0.002)	0.654 (0.001)
145	13-1157	13-490	−0.287 (0.003)	6504.274 (0.004)	13.587 (0.049)	140.112 (39.911)	0.116 (0.065)	−1.493 (0.581)	15.260 (0.000)	0.005 (0.001)
146	13-1227	13-545	−0.052 (0.003)	6515.517 (0.027)	93.445 (3.798)	13.861 (10.893)	−0.165 (0.022)	0.803 (0.121)	17.576 (0.001)	0.924 (0.000)
147	13-1250	13-516	0.552 (0.117)	6491.580 (0.019)	2.588 (0.302)	582.467 (208.685)	−0.264 (0.801)	0.240 (0.789)	15.668 (0.000)	0.925 (0.001)
148	13-1253	13-519	−0.110 (0.006)	6497.594 (0.008)	7.984 (0.260)	18.925 (27.287)	1.469 (0.593)	1.476 (0.803)	17.138 (0.001)	0.786 (0.001)
149	13-1257	13-556	0.253 (0.068)	6514.259 (0.038)	23.877 (0.839)	128.485 (85.036)	0.994 (0.157)	−1.475 (0.771)	18.625 (0.001)	0.513 (0.002)
150	13-1279	13-518	0.082 (0.001)	6499.373 (0.001)	3.808 (0.010)	<b>139.541</b> ( <b>0.859</b> )	0.000 (0.000)	0.000 (0.000)	15.799 (0.000)	0.000 (0.000)
151	13-1287	13-537	0.576 (0.051)	6503.415 (0.046)	13.465 (0.710)	0.000 (155.043)	0.163 (0.314)	−1.446 (0.839)	17.502 (0.001)	0.314 (0.003)
152	13-1492	13-548	−0.243 (0.074)	6504.597 (0.054)	13.462 (2.230)	254.502 (111.644)	−1.064 (0.757)	0.286 (0.795)	17.274 (0.002)	0.911 (0.002)
153	13-1498	13-568	−0.232 (0.010)	6517.917 (0.031)	18.762 (0.403)	184.671 (59.591)	0.737 (0.178)	−1.410 (0.873)	18.842 (0.002)	0.000 (0.000)
154	13-1506	13-541	−0.110 (0.004)	6512.931 (0.002)	12.343 (0.125)	105.627 (8.549)	−0.742 (0.135)	0.612 (0.816)	16.009 (0.001)	0.017 (0.002)
155	13-1551	13-572	−0.155 (0.022)	6517.452 (0.024)	15.373 (1.063)	163.872 (53.582)	−0.963 (0.505)	−0.276 (0.833)	19.404 (0.005)	0.506 (0.005)
156	13-1598	13-570	−0.754 (0.035)	6532.931 (0.024)	10.952 (0.325)	4.218 (69.121)	<b>−1.500</b> ( <b>0.019</b> )	<b>−1.461</b> ( <b>0.858</b> )	15.584 (0.000)	0.133 (0.003)
157	13-1617	13-560	−0.200 (0.007)	6515.385 (0.007)	11.019 (0.225)	163.536 (41.817)	−0.865 (0.296)	−0.955 (0.777)	17.642 (0.001)	0.152 (0.003)
158	13-1669	13-585	−0.007 (0.000)	6527.163 (0.003)	52.476 (2.138)	9.301 (0.471)	0.163 (0.111)	1.496 (1.054)	20.463 (0.019)	0.013 (0.019)
159	13-1678	13-595	−0.204 (0.028)	6531.602 (0.018)	7.545 (0.479)	0.000 (77.941)	1.494 (0.469)	1.379 (0.811)	18.312 (0.003)	0.413 (0.006)
160	13-1689	13-593	−0.081 (0.003)	6534.384 (0.001)	5.177 (0.062)	47.757 (20.507)	0.931 (0.646)	−0.432 (0.794)	17.420 (0.001)	0.086 (0.003)

Table A1 – continued

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E,E}$	$\pi_{E,N}$	$I_{bl}$	$f_{bl}$
161	13-1721	13-618	−0.086 (0.009)	6535.020 (0.042)	22.085 (2.239)	92.966 (26.987)	−0.336 (0.447)	−1.119 (0.825)	20.502 (0.010)	0.254 (0.011)
162	13-1730	13-602	−0.214 (0.021)	6538.838 (0.041)	17.116 (0.940)	3.528 (64.806)	−0.216 (0.414)	−0.140 (0.803)	18.848 (0.004)	0.330 (0.007)
163	13-1738	13-628	0.286 (0.075)	6542.673 (0.042)	12.642 (0.809)	8.011 (87.152)	−1.265 (0.427)	0.079 (0.776)	18.951 (0.002)	0.328 (0.005)
164	13-1775	13-614	−0.008 (0.001)	6538.850 (0.004)	46.183 (7.734)	7.671 (1.873)	−0.060 (0.300)	−0.392 (0.750)	20.126 (0.012)	0.750 (0.003)
165	14-0099	14-109	−0.343 (0.005)	6876.998 (0.087)	123.255 (1.073)	263.093 (71.902)	<b>0.194</b> <b>(0.004)</b>	<b>0.291</b> <b>(0.020)</b>	16.683 (0.000)	0.338 (0.001)
166	14-0115	14-273	−0.348 (0.009)	6859.582 (0.040)	97.229 (0.858)	307.209 (21.747)	<b>0.110</b> <b>(0.005)</b>	<b>0.039</b> <b>(0.030)</b>	17.117 (0.001)	0.010 (0.001)
167	14-0124	14-307	−0.160 (0.004)	6836.531 (0.047)	167.579 (5.816)	89.568 (38.330)	0.129 (0.014)	−0.074 (0.027)	17.382 (0.001)	0.712 (0.001)
168	14-0257	14-148	0.001 (0.006)	6772.704 (0.028)	61.720 (1.309)	139.227 (2.858)	−0.140 (0.015)	−0.078 (0.153)	19.027 (0.003)	0.333 (0.003)
169	14-0337	14-194	−0.543 (0.006)	6822.955 (0.023)	45.276 (0.187)	207.333 (78.241)	<b>0.146</b> <b>(0.042)</b>	<b>1.416</b> <b>(0.075)</b>	16.768 (0.000)	0.000 (0.000)
170	14-0383	14-147	0.013 (0.004)	6761.393 (0.001)	15.139 (0.293)	10.691 (3.899)	−0.041 (0.171)	−0.549 (0.753)	18.032 (0.002)	0.583 (0.001)
171	14-0419	14-283	−0.248 (0.012)	6822.954 (0.024)	47.010 (1.166)	0.000 (57.567)	−0.551 (0.175)	−0.861 (0.333)	17.918 (0.002)	0.194 (0.002)
172	14-0512	14-157	0.675 (0.003)	6782.126 (0.010)	11.354 (0.038)	0.000 (40.879)	<b>0.914</b> <b>(0.078)</b>	<b>1.492</b> <b>(0.259)</b>	15.679 (0.000)	0.000 (0.000)
173	14-0530	14-220	−0.117 (0.002)	6791.737 (0.005)	20.530 (0.162)	51.944 (18.433)	−0.034 (0.103)	−0.372 (0.650)	17.684 (0.001)	0.093 (0.001)
174	14-0531	14-229	0.186 (0.029)	6784.771 (0.071)	41.753 (3.073)	1.559 (56.852)	0.151 (0.173)	−1.499 (0.792)	18.045 (0.001)	0.802 (0.001)
175	14-0572	14-165	−0.561 (0.007)	6775.903 (0.036)	9.009 (0.083)	3.733 (65.306)	1.498 (0.052)	−1.480 (0.555)	17.445 (0.001)	0.000 (0.000)
176	14-0582	14-151	0.005 (0.028)	6766.182 (0.026)	2.074 (0.245)	84.547 (32.638)	−1.127 (0.793)	0.463 (0.805)	19.220 (0.004)	0.453 (0.006)
177	14-0617	14-215	0.141 (0.085)	6771.464 (0.007)	1.565 (0.162)	236.811 (67.919)	0.469 (0.802)	−0.516 (0.778)	18.756 (0.003)	0.718 (0.004)
178	14-0676	14-175	0.004 (0.077)	6778.228 (0.226)	8.015 (5.435)	110.012 (50.335)	−0.910 (0.828)	0.487 (0.789)	19.537 (0.016)	0.456 (0.026)
179	14-0692	14-214	0.255 (0.057)	6778.824 (0.037)	7.763 (0.737)	274.031 (74.266)	1.199 (0.738)	−0.476 (0.769)	18.853 (0.002)	0.762 (0.004)
180	14-0696	14-199	0.405 (0.001)	6788.774 (0.003)	6.832 (0.011)	3.167 (24.469)	<b>−1.393</b> <b>(0.098)</b>	<b>1.496</b> <b>(0.512)</b>	14.183 (0.000)	0.000 (0.000)
181	14-0734	14-243	−0.753 (0.046)	6791.432 (0.027)	5.851 (0.240)	424.005 (176.220)	1.483 (0.707)	−0.740 (0.813)	17.725 (0.001)	0.019 (0.005)
182	14-0727	14-281	1.043 (0.162)	6842.546 (2.346)	38.110 (4.552)	0.000 (306.481)	−0.819 (0.142)	−0.075 (0.127)	16.607 (0.001)	0.056 (0.013)
183	14-0729	14-262	−0.663 (0.008)	6799.616 (0.017)	6.557 (0.045)	4.023 (85.907)	−1.226 (0.486)	1.430 (0.778)	16.463 (0.000)	0.000 (0.000)
184	14-0783	14-253	−0.017 (0.001)	6791.411 (0.002)	21.003 (1.410)	<b>19.919</b> <b>(1.993)</b>	<b>0.072</b> <b>(0.488)</b>	<b>1.187</b> <b>(0.794)</b>	20.117 (0.005)	0.532 (0.003)
185	14-0874	14-302	−0.199 (0.004)	6845.679 (0.006)	24.877 (0.096)	130.938 (41.271)	−0.030 (0.083)	1.482 (0.351)	15.887 (0.000)	0.000 (0.000)
186	14-0893	14-296	−0.024 (0.007)	6809.366 (0.014)	74.901 (8.718)	0.000 (6.110)	0.000 (0.243)	0.000 (0.551)	19.660 (0.007)	0.857 (0.002)
187	14-0894	14-333	−0.458 (0.010)	6846.403 (0.026)	28.577 (0.491)	0.000 (71.085)	0.213 (0.074)	−1.499 (0.870)	17.530 (0.001)	0.031 (0.003)
188	14-0900	14-305	0.280 (0.010)	6862.267 (0.041)	33.421 (0.238)	0.000 (73.755)	−0.202 (0.022)	−1.499 (0.017)	16.060 (0.001)	0.000 (0.000)
189	14-0928	14-361	−0.387 (0.014)	6877.343 (0.043)	55.656 (1.576)	106.812 (66.610)	−0.076 (0.023)	−1.402 (0.068)	16.570 (0.000)	0.181 (0.001)
190	14-9958	14-275	0.090 (0.001)	6817.268 (0.031)	79.072 (0.290)	34.569 (13.950)	0.120 (0.022)	0.884 (0.049)	19.074 (0.002)	0.000 (0.000)
191	14-0962	14-285	−0.002 (0.027)	6817.988 (0.014)	6.540 (0.085)	199.230 (10.398)	−1.476 (0.237)	1.493 (0.622)	16.171 (0.000)	0.524 (0.002)
192	14-0996	14-310	−0.064 (0.002)	6812.742 (0.007)	16.899 (0.150)	1.083 (15.487)	1.257 (0.650)	1.409 (0.812)	20.275 (0.003)	0.000 (0.000)



Table A1 – *continued*

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E, E}$	$\pi_{E, N}$	$I_{bl}$	$f_{bl}$
193	14-1029	14-312	0.002 (0.001)	6811.241 (0.005)	125.789 (53.289)	0.950 (0.855)	−1.425 (0.282)	0.101 (0.381)	20.353 (0.006)	0.984 (0.000)
194	14-1101	14-362	0.713 (0.035)	6855.431 (0.045)	20.393 (0.499)	358.065 (167.500)	−0.897 (0.319)	1.372 (0.745)	16.892 (0.000)	0.000 (0.000)
195	14-1109	14-328	0.785 (0.085)	6826.141 (0.076)	7.855 (0.681)	608.533 (205.985)	−1.414 (0.805)	1.499 (0.747)	18.304 (0.002)	0.044 (0.011)
196	14-1111	14-401	−0.181 (0.005)	6861.902 (0.026)	32.831 (0.587)	2.518 (36.410)	0.077 (0.089)	0.676 (0.706)	18.740 (0.003)	0.006 (0.004)
197	14-1153	14-389	−0.366 (0.008)	6853.536 (0.018)	16.342 (0.174)	0.000 (68.999)	0.081 (0.216)	−1.485 (0.837)	17.928 (0.001)	0.001 (0.002)
198	14-1255	14-367	0.011 (0.013)	6844.289 (0.006)	25.208 (1.188)	0.000 (11.132)	0.050 (0.589)	1.400 (0.770)	18.551 (0.004)	0.636 (0.002)
199	14-1284	14-392	0.037 (0.009)	6852.991 (0.010)	40.898 (1.745)	0.000 (7.153)	0.551 (0.217)	−0.992 (0.785)	19.230 (0.003)	0.687 (0.001)
200	14-1285	14-370	−0.567 (0.007)	6848.987 (0.008)	6.316 (0.033)	0.000 (72.928)	1.500 (0.464)	1.180 (0.812)	16.076 (0.000)	0.000 (0.000)
201	14-1308	14-357	0.493 (0.103)	6846.555 (0.040)	10.199 (0.233)	0.000 (142.274)	−0.894 (0.707)	−1.372 (0.792)	17.794 (0.003)	0.061 (0.014)
202	14-1317	14-364	−0.045 (0.007)	6840.933 (0.032)	40.873 (4.059)	0.000 (16.625)	−1.484 (0.708)	−0.200 (0.744)	20.254 (0.014)	0.651 (0.006)
203	14-1340	14-372	−0.084 (0.003)	6845.154 (0.004)	6.697 (0.173)	0.085 (18.720)	1.491 (0.811)	1.497 (0.819)	19.393 (0.004)	0.000 (0.000)
204	14-1392	14-398	0.097 (0.007)	6854.863 (0.010)	13.448 (0.571)	97.901 (22.785)	−0.434 (0.593)	−0.910 (0.795)	20.017 (0.004)	0.191 (0.005)
205	14-1406	14-409	−0.344 (0.003)	6873.449 (0.026)	12.544 (0.065)	0.000 (49.425)	0.100 (0.195)	−1.497 (0.827)	15.840 (0.001)	0.000 (0.000)
206	14-1409	14-393	−0.029 (0.006)	6853.793 (0.010)	11.255 (1.519)	14.686 (12.319)	−1.192 (0.794)	0.144 (0.810)	18.945 (0.002)	0.890 (0.001)
207	14-1410	14-424	0.199 (0.023)	6868.477 (0.027)	21.526 (0.883)	195.626 (71.835)	−0.268 (0.230)	−1.486 (0.844)	18.435 (0.002)	0.434 (0.004)
208	14-1437	14-407	−0.056 (0.020)	6858.442 (0.006)	7.216 (0.354)	0.000 (15.973)	1.463 (0.819)	1.132 (0.793)	18.466 (0.005)	0.426 (0.004)
209	14-1472	14-433	0.147 (0.006)	6864.195 (0.022)	12.511 (0.413)	40.003 (34.658)	0.426 (0.705)	0.926 (0.814)	19.705 (0.006)	0.007 (0.009)
210	14-1483	14-506	0.063 (0.003)	6895.944 (0.029)	54.295 (2.656)	57.645 (14.109)	−0.009 (0.060)	0.261 (0.394)	18.862 (0.002)	0.762 (0.001)
211	14-1486	14-423	−0.107 (0.004)	6864.709 (0.001)	5.525 (0.055)	82.538 (26.708)	1.428 (0.535)	1.499 (0.813)	16.596 (0.000)	0.565 (0.000)
212	14-1498	14-456	−0.150 (0.006)	6881.422 (0.032)	25.862 (0.753)	0.000 (33.263)	0.000 (0.112)	0.037 (0.659)	19.411 (0.002)	0.336 (0.003)
213	14-1501	14-426	0.007 (0.001)	6859.360 (0.002)	43.698 (2.755)	0.000 (2.010)	0.004 (0.215)	0.093 (0.701)	20.328 (0.012)	0.099 (0.011)
214	14-1570	14-451	−0.697 (0.025)	6887.094 (0.017)	12.102 (0.303)	451.023 (130.543)	0.236 (0.110)	0.752 (0.737)	16.121 (0.000)	0.169 (0.001)
215	14-1812	14-444	0.017 (0.002)	6873.149 (0.017)	68.707 (14.538)	26.821 (3.729)	<b>1.459</b> (0.187)	<b>1.003</b> (0.797)	21.368 (0.072)	0.243 (0.053)
216	14-1596	14-445	−0.206 (0.014)	6874.945 (0.057)	16.352 (0.760)	210.268 (53.485)	−1.476 (0.343)	−0.359 (0.771)	18.457 (0.003)	0.575 (0.005)
217	14-1598	14-448	−0.096 (0.031)	6874.447 (0.008)	3.484 (0.238)	0.844 (28.613)	−1.488 (0.823)	−1.155 (0.801)	19.253 (0.004)	0.121 (0.008)
218	14-1600	14-494	−0.285 (0.039)	6889.198 (0.058)	27.047 (1.592)	0.000 (113.552)	−0.198 (0.115)	1.486 (0.812)	19.070 (0.003)	0.277 (0.005)
219	14-1647	14-489	0.164 (0.012)	6892.046 (0.018)	16.237 (0.580)	0.000 (49.741)	−0.086 (0.310)	−1.482 (0.842)	17.716 (0.001)	0.840 (0.001)
220	14-1701	14-488	−0.124 (0.005)	6891.664 (0.005)	6.797 (0.162)	122.430 (9.349)	−1.459 (0.652)	−1.397 (0.790)	17.559 (0.001)	0.668 (0.001)
221	14-1706	14-502	−0.807 (0.038)	6892.816 (0.024)	3.298 (0.104)	0.000 (141.679)	−1.500 (0.729)	−0.994 (0.815)	17.221 (0.001)	0.000 (0.010)
222	14-1720	14-507	−0.045 (0.004)	6894.912 (0.006)	13.556 (0.817)	75.223 (5.115)	−1.370 (0.631)	−0.662 (0.788)	19.072 (0.005)	0.718 (0.004)
223	14-1748	14-496	−0.001 (0.002)	6891.179 (0.001)	6.973 (0.976)	0.000 (1.931)	−1.349 (0.796)	−1.282 (0.834)	19.827 (0.011)	0.664 (0.006)
224	14-1816	14-498	0.065 (0.003)	6893.400 (0.008)	13.480 (0.293)	18.307 (15.881)	−1.496 (0.140)	−1.411 (0.765)	19.947 (0.005)	0.000 (0.000)

Table A1 – continued

#	OGLE no.	MOA no.	$u_0$	$t_0 - 2450000$ (HJD)	$t_E$ (d)	$\rho$ ( $10^{-3}$ )	$\pi_{E,E}$	$\pi_{E,N}$	$I_{bl}$	$f_{bl}$
X1	11-0975	11-305	0.408 (0.252)	5760.427 (0.238)	14.646 (2.751)	1145.523 (162.081)	−0.473 (0.738)	0.120 (0.795)	19.344 (0.007)	0.009 (0.043)
X2	12-0937	12-405	−0.063 (0.131)	6102.288 (0.099)	9.248 (0.523)	725.970 (37.411)	1.226 (0.806)	1.221 (0.800)	20.328 (0.008)	0.071 (0.019)
X3	12-1502	12-567	0.064 (0.041)	6138.430 (2.725)	393.584 (79.223)	71.659 (80.405)	−0.810 (0.310)	0.053 (0.229)	20.365 (0.012)	0.933 (0.002)
X4	13-0966	13-428	−0.352 (0.236)	6450.141 (0.143)	8.388 (11.326)	223.542 (221.974)	1.180 (0.796)	−0.863 (0.788)	20.108 (0.013)	0.404 (0.026)
X5	14-0289	14-092	0.572 (0.170)	6852.483 (28.456)	230.254 (113.809)	837.450 (395.201)	−0.260 (0.456)	1.061 (0.200)	18.813 (0.031)	0.243 (0.082)
X6	14-0921	14-270	−0.003 (0.234)	6803.854 (0.145)	374.898 (90.046)	3.852 (151.397)	0.490 (0.788)	0.229 (0.819)	19.724 (0.021)	0.996 (0.000)

Note. Significant detections of finite source effects and parallax are marked in bold.

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